

## The Phase function in the Tertiary Scattering

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### Abstract

Let us take at random three points F, T and E in the earth's atmosphere, and a point O' on its surface. Define a system of rectangular coordinates  $X_1, Y_1, Z_1$  with its center at F, with  $X_1$  axis drawn towards the Sun. Let us resolve the direct insolation reaching F into two plane polarized rays: the one travels to X<sub>1</sub> and oscillates in Z<sub>1</sub> direction, which is named by (1), the other travels to X<sub>1</sub> and oscillates in Y<sub>1</sub>, which is named by (2). The primarily scattered ray generated at T when (1) encounters one air particle at T travels to FT direction and oscillates in a direction normal to it. FT direction is named by X<sub>2</sub>. This direction of oscillation is named by Z<sub>2</sub> and determined by X<sub>1</sub>, Y<sub>1</sub>, Z<sub>1</sub> and the positions of F and T. This scattering is named by E<sub>1</sub>.

The primarily scattered ray generated at T when (2) encounters one air particle at T travels to FT direction and oscillates in a direction normal to FT. This direction is named by Z'<sub>2</sub> and determined by X<sub>1</sub>, Y<sub>1</sub>, Z<sub>1</sub> and F, T. This scattering is named by E'<sub>1</sub>.

E<sub>1</sub> travels to TE direction and oscillates in a direction normal to it at the point T. This direction of oscillation is named by Z<sub>3</sub>. E<sub>1</sub> scattering generates a secondarily scattered ray at E when it encounters one air particle at E. This secondary scattering is named by E<sub>2</sub>. In the same way, E'<sub>1</sub> travels to TE direction and oscillates in a direction normal to it at the point T. This direction of oscillation is named by Z'<sub>3</sub>. E'<sub>1</sub> scattering generates a secondarily scattered ray at E when it encounters one air particle at E. This secondary scattering is named by E'<sub>2</sub>.

Further, E<sub>2</sub> and E'<sub>2</sub> will generate tertiary scattering at O' when they reach at O' and encounter one air particle at O'.

Let  $\omega_1, \omega_1', \omega_2, \omega_2', \omega_3, \omega_3'$  be respectively the angle between FT and Z<sub>1</sub>, FT and Y<sub>1</sub>, TE and Z<sub>2</sub>, TE and Z'<sub>2</sub>, EO' and Z<sub>3</sub>, EO' and Z'<sub>3</sub>, then the phase function in the tertiary scattering is expressed by

$$\sum_{n=1}^3 \pi \sin^2 \omega_n + \sum_{n=1}^3 \pi \sin^2 \omega_n'$$

Let  $O$  be the earth's center. Define a system of rectangular coordinate  $X' Y' Z'$  with its center at  $O'$ .  $Z$  is drawn towards  $OO'$ ,  $X'$  is in the vertical plane containing the Sun's centre and in the Sun's side. Let  $E, T$  and  $F$  be in the plane  $X' Z'$  in the Sun's side and use the notation  $\angle E O' X' = \theta_1$ ,  $\angle T E O = \theta_1$ ,  $\angle E O T = \theta_2$ ,  $\angle O T F = \theta_3$ ,  $h$  = the Sun's altitude. Then the phase function are calculated as follows in the case when  $E$  approaches infinitely to  $O'$  with a definite value of  $\theta$ .  $D+D' = 1 + \vartheta \sin^2(\theta + \theta_1)$ , here  $\vartheta$  being a function of

$$\theta_1, \theta_2, \theta_3, h.$$

1. Introduction The scattering phenomena of the Sun's ray by the earth's atmosphere is the original cause of the blue sky in the day time and the yellowish and reddish sky in the morning and evening. In the scattering problem the terms scattering and absorption are in general used separately as if they were different from each other, but the latter is originated from the former in the Rayleigh atmosphere. Hence we can reasonably say that the color of the sky is caused only by scattering with no suspension. The scattering phenomena has been solved at first by Rayleigh by the consideration of dimension (Ref. 1). Then, Max Planck has established the theory of electromagnetic wave and made the theoretical formula of the primary scattering intensity basing on the field equation (Ref. 2).

The author has expanded his research and established the theoretical formula giving the phase function in the secondary scattering and the method of computation and moreover executed the computation precisely (Ref. 3).

2. Theoretical formula Take three points  $F, T$  and  $E$  in the atmosphere and a point  $O'$  on the earth's surface (see fig. 1). Define a system of rectangular coordinates  $X_1, Y_1, Z_1$ , with its centre at  $F$ ,  $X_1$  axis drawn towards the Sun.  $Z_1$  axis is in the vertical plane passing through the Sun's centre at  $O'$ . Let  $(r_1, \delta_1, \kappa_1)$  be the coordinates of  $T$  of this system. Now, if  $i_1$  be the direct insolation reaching  $F$ , we can resolve it into two plane polarized rays, the one is

$$E(Z_1) = \sqrt{\frac{i_1}{2}} \exp(int + imx) \quad (1)$$

which travels to  $X_1$  and oscillates in  $Z_1$  direction, the other is

$$E(Y_1) = \sqrt{\frac{i_1}{2}} \exp(int + imx) \quad (2)$$

travelling to  $X_1$  and oscillating in  $Y_1$  direction.

Here  $FT = R_1$  and  $\epsilon$  be the dielectric constant of air particles. Then the equation of wave generated at  $T$  when the former ie  $E(Z_1)$  encounters one air particle at  $T$  is expressed by

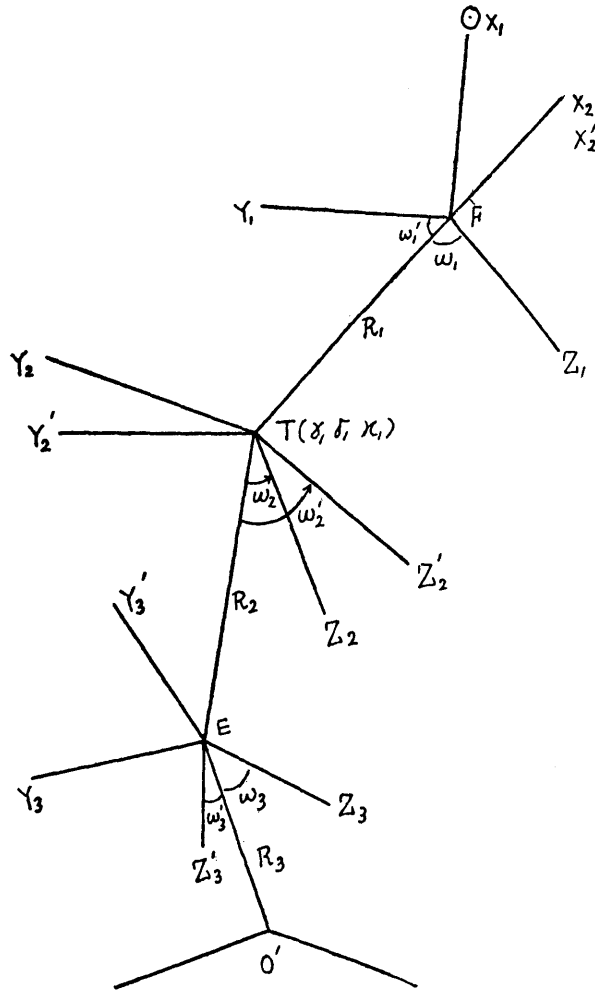


Fig. 1 Explanation of the polarization angle in the tertiary scattering.

$$\left. \begin{aligned}
 E_{ix} &= -\sqrt{\frac{i_\lambda}{2}} \frac{\pi T}{R_1 \lambda^2} e^{(\Delta \frac{1}{\epsilon})} \frac{\gamma_1 \kappa_1}{R_1^2} \exp(int - imR_1) \\
 E_{iy} &= -\sqrt{\frac{i_\lambda}{2}} \frac{\pi T}{R_1 \lambda^2} e^{(\Delta \frac{1}{\epsilon})} \frac{\delta_1 \kappa_1}{R_1^2} \exp(int - imR_1) \\
 E_{iz} &= -\sqrt{\frac{i_\lambda}{2}} \frac{\pi T}{R_1 \lambda^2} e^{(\Delta \frac{1}{\epsilon})} \frac{\gamma_1^2 + \delta_1^2}{R_1^2} \exp(int - imR_1)
 \end{aligned} \right\} (3)$$

in which  $\lambda$  and  $T$  are the wave length and the volume of particle. This wave will proceed into the direction  $FT$  and have a direction of oscillation normal to it, and the direction cosine of the direction of oscillation are as follows :

$$\frac{1}{\sin \omega_1} \frac{\gamma_1 \kappa_1}{R_1^2}, \frac{1}{\sin \omega_1} \frac{\delta_1 \kappa_1}{R_1^2}, -\frac{1}{\sin \omega_1} \frac{\gamma_1^2 + \delta_1^2}{R_1^2} \quad (4)$$

in which  $\sin \omega_1 = \frac{\sqrt{\gamma_1^2 + \delta_1^2}}{R_1}$ ,

i. e.  $\omega_1$  is the angle between FT and  $Z_1$  axis. The intensity can be obtained by squaring the amplitude, which is

$$E_1 = \sqrt{\frac{i_1}{2}} \frac{\pi T}{R_1 \lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right) \sin \omega_1 \quad (5)$$

In the same way, the plane polarized light generated at T by  $E(Y_1)$  is

$$\left. \begin{aligned} E'_{1x} &= \sqrt{\frac{i_1}{2}} \frac{\pi T}{R_1 \lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right) \frac{\gamma_1 \delta_1}{R_1^2} \exp(\text{int} - imR_1) \\ E'_{1y} &= -\sqrt{\frac{i_1}{2}} \frac{\pi T}{R_1 \lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right) \frac{\gamma_1^2 + \kappa_1^2}{R_1^2} \exp(\text{int} - imR_1) \\ E'_{1z} &= \sqrt{\frac{i_1}{2}} \frac{\pi T}{R_1 \lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right) \frac{\delta_1 \kappa_1}{R_1^2} \exp(\text{int} - imR_1) \end{aligned} \right\} (6)$$

This will also proceed into FT direction and oscillate normal to it, the amplitude being

$$E'_1 = \sqrt{\frac{i_1}{2}} \frac{\pi T}{R_1 \lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right) \sin \omega'_1 \quad (7)$$

in which  $\omega'_1$  is the angle between FT and  $Y_1$  axis. In the following discussion, let the amplitude  $E_1$  and  $E'_1$  also be the notation of their corresponding polarizations.

Now take the origin at T,  $X_2$  axis in TF direction,  $Z_2$  axis in the direction of oscillation of  $E_1$ ,  $Y_2$  axis normal to  $X_2 Z_2$ . In the same way, define  $X'_2 Y'_2 Z'_2$  with respect to  $E'_1$ , in which  $X'_2$  being identical with  $X_2$ .

Now, take a point E, and let  $TE = R_2$  and the angular distances of TE from  $Z_2$  and  $Z'_2$  axes be  $\omega_2$  and  $\omega'_2$ . Then,  $E_1$  at T will also generate one plane polarized light at E. Letting the coordinates of E referred to  $X_2 Y_2 Z_2$  system be  $(r_2 \delta_2 \kappa_2)$ , then this new plane polarized light will be

$$\left. \begin{aligned} E_{2x} &= E_1 \frac{\pi T}{R_2 \lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right) \frac{\gamma_2 \kappa_2}{R_2^2} \exp(\text{int} - imR_2) \\ E_{2y} &= E_1 \frac{\pi T}{R_2 \lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right) \frac{\delta_2 \kappa_2}{R_2^2} \exp(\text{int} - imR_2) \\ E_{2z} &= E_1 \frac{\pi T}{R_2 \lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right) \frac{\gamma_2^2 + \delta_2^2}{R_2^2} \exp(\text{int} - imR_2) \end{aligned} \right\} (8)$$

whose amplitude being

$$E_2 = E_1 \frac{\pi T}{R_2 \lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right) \sin \omega_2 \quad (9)$$

In the same way, the plane polarized light generated by  $E'_1$  at E point can be

expressed by substituting  $E_2'E_1' \gamma_2'\delta_2'\kappa_2'\omega_2'$  for  $E_2E_1 \gamma_2\delta_2\kappa_2\omega_2$  in the above expressions with respect to  $X_2'Y_2'Z_2'$  system. Here, both directions of oscillation of  $E_2$  and  $E_2'$  are naturally normal to TE. Now, taking the origin at E,  $X_3$  axis in ET direction,  $Z_3'$  in the direction of oscillation of  $E_2$ ,  $Y_3$  normal to them,  $X_3'$  in ET direction,  $Z_3'$  in the direction of oscillation of  $E_2$ ,  $Y_3'$  normal to them, and letting the coordinates of  $O'$  with respect to these systems be  $\gamma_3\delta_3\kappa_3$ ,  $\gamma_3'\delta_3'\kappa_3'$ ,  $EO'=R_3$  and the angular distances of  $EO'$  from  $Z_3$  and  $Z_3'$  be  $\omega_3$  and  $\omega_3'$ , then the plane polarized light generated at  $O'$  by  $E_2$  is expressed by referring  $X_3 Y_3 Z_3$

$$E_{3x} = E_2 \frac{\pi T}{R_3 \lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right) \frac{\gamma_3 \kappa_3}{R_3^2} \exp(\text{int} - \text{im}R_3)$$

$$E_{3y} = E_2 \frac{\pi T}{R_3 \lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right) \frac{\delta_3 \kappa_3}{R_3^2} \exp(\text{int} - \text{im}R_3) \quad (10)$$

$$E_{3z} = -E_2 \frac{\pi T}{R_3 \lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right) \frac{\gamma_3^2 + \delta_3^2}{R_3^2} \exp(\text{int} - \text{im}R_3)$$

whose amplitude being

$$E_3 = E_2 \frac{\pi T}{R_3 \lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right) \sin \omega_3 \quad (11)$$

In the same way, the plane polarized light generated by  $E_2'$  at  $O'$  point can be expressed by substituting  $E_2' \gamma_3' \delta_3' \kappa_3' \omega_3'$  for  $E_2 \gamma_3 \delta_3 \kappa_3 \omega_3$  in the above expressions with respect to  $X_3'Y_3'Z_3'$  system.

Eventually we can get

$$E_3 = \sqrt{\frac{i_\lambda}{2}} \left(\frac{\pi T}{\lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right)\right)^3 \frac{\sin \omega_1}{R_1} \cdot \frac{\sin \omega_2}{R_2} \cdot \frac{\sin \omega_3}{R_3}, \quad (12)$$

$$E_3' = \sqrt{\frac{i_\lambda}{2}} \left(\frac{\pi T}{\lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right)\right)^3 \frac{\sin \omega_1'}{R_1} \cdot \frac{\sin \omega_2'}{R_2} \cdot \frac{\sin \omega_3'}{R_3}$$

Now, let the number of particles in unit volume at F, T and E be  $I_F$ ,  $I_T$  and  $I_E$  respectively and put

$$k_{in} = \frac{8 \pi^3 T^3}{3 \lambda^4} I_n \epsilon^3 \left(\Delta \frac{1}{\epsilon}\right)^3 \quad (13)$$

(n = F, T, E)

As the energy is the square of the amplitude, so the intensity of tertiary scattering generated at  $O'$  by unit volume of F, T and E will be

$$\frac{i_\lambda}{2} \left(\frac{\pi T}{\lambda^2} \epsilon \left(\Delta \frac{1}{\epsilon}\right)\right)^6 \frac{I_F}{R_1^2} \cdot \frac{I_T}{R_2^2} \cdot \frac{I_E}{R_3^2} (\sin^2 \omega_1 \sin^2 \omega_2 \sin^2 \omega_3 + \sin^2 \omega_1' \sin^2 \omega_2' \sin^2 \omega_3')$$

$$= \frac{i_\lambda}{2} \left(\frac{3}{8 \pi}\right)^3 \frac{1}{R_1^2} \cdot \frac{1}{R_2^2} \cdot \frac{1}{R_3^2} \left(\sum_{n=1}^3 \sin^2 \omega_n + \sum_{n=1}^3 \sin^2 \omega_n'\right) k_{xF} k_{xT} k_{xE} \quad (14)$$

If we consider the effect of absorption on the optical way, we must only

multiply the absorption terms.

3. Practical method of calculation. Let  $O$  be the earth's centre and  $O'$  a point on its surface. Take a coordinate system  $X Y Z$  with its origin at  $O$ ,  $Z$  axis being directed towards  $OO'$ ,  $X$  axis normal to  $Z$  axis and towards the Sun's side on the plane containing  $OO'$  and its centre,  $Y$  axis normal to  $X$  and  $Z$ .

Take a point  $E$  in the atmosphere seen from  $O'$  and let the coordinates referred to  $X Y Z$  system be

$$X = e \sin r \cos A, \quad Y = e \sin r \sin A, \quad Z = e \cos r \quad (15)$$

in which

$$0 \leq r \leq \frac{\pi}{2} \quad 0 \leq A \leq 2\pi$$

Let  $X' Y' Z'$  system be the parallel translation of  $X Y Z$  system by the transformation of the origin  $O$  to  $O'$ , and put (see fig. 2)

$$X' = R \cos \theta \cos A, \quad Y' = R \cos \theta \sin A, \quad Z' = R \sin \theta \quad (16)$$

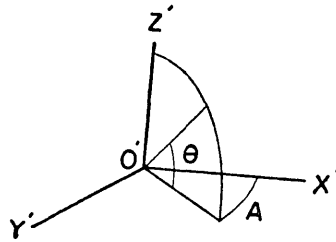


Fig. 2  $X'Y'Z'$  coordinate system and angles  $A, \theta$ .

Letting the coordinates of  $E$  referred to this new system be  $X' Y' Z'$ , so we get

$$X = X', \quad Y = Y', \quad Z = Z' + a^\circ \quad (17)$$

where  $a^\circ$  being the earth's radius.

The relation of  $(x_2 y_2 z_2)$  coordinate system to  $(X Y Z)$  system is defined as follows

	$x_2$	$y_2$	$z_2$	
X	$\cos r \cos A$	$-\sin A$	$\sin r \cos A$	(18)
Y	$\cos r \sin A$	$\cos A$	$\sin r \sin A$	
Z	$-\sin r$	0	$\cos r$	

Now let the coordinates of  $T$  referred to  $(x_2 y_2 z_2)$  system be

$$x_{2T} = OT \sin \theta \cos A, \quad y_{2T} = OT \sin \theta \sin A, \quad z_{2T} = OT \cos \theta \quad (19)$$

(see fig. 3)

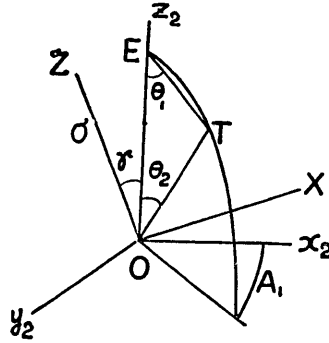


Fig. 3 Relation of  $x_2 y_2 z_2$  system to  $X Y Z$  and definition of angles  $r, \theta_1, \theta_2, A_1$ .

Let  $(x_2' y_2' z_2')$  system be the parallel transformation of  $(x_2 y_2 z_2)$  system by the translation of the origin  $O$  to  $E$ , and put  $\angle EOT = \theta_2, \angle OET = \theta_1$ , then

$$\sin \theta_2 = \frac{ET}{OT} \sin \theta_1 \tag{20}$$

and the coordinates of  $T$  referred to  $(x_2' y_2' z_2')$  will be

$$x_2' = ET \sin \theta_1 \cos A_1, \quad y_2' = ET \sin \theta_1 \sin A_1, \quad z_2' = -ET \cos \theta_1 \tag{21}$$

Giving  $ET, \theta_1, A_1$ , we can get  $OT$ , so  $\theta_2$  can be found, and eventually  $x_{2T}, y_{2T}, z_{2T}$  i. e. the coordinate of  $T$  referring to  $(x_2 y_2 z_2)$  system.

Moreover, let the polar coordinates of  $T$  referred to  $(X Y Z)$  system be

$$X_T = OT \sin r' \cos A', \quad Y_T = OT \sin r' \sin A', \quad Z_T = OT \cos r' \tag{22}$$

(see fig. 4)

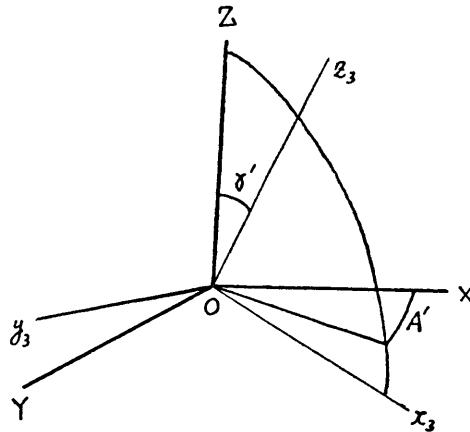


Fig. 4 Relation of  $x_3 y_3 z_3$  system to  $XYZ$  system and definition of angles  $A', r'$ .

The combination of the relation between  $X Y Z$  and  $x_2 y_2 z_2$  systems (18) and the polar coordinates of  $x_2 y_2 z_2$  system will generate the following formula

$$\begin{aligned} \sin r' \cos A' &= \cos r \cos A \sin \theta_2 \cos A_1 - \sin A \sin \theta_2 \sin A_1 + \sin r \cos A \cos \theta_2 \\ \sin r' \sin A' &= \cos r \sin A \sin \theta_2 \cos A_1 + \cos A \sin \theta_2 \sin A_1 + \sin r \sin A \cos \theta_2 \\ \cos r' &= -\sin r \sin \theta_2 \cos A_1 + \cos r \cos \theta_2 \end{aligned} \quad (23)$$

From this we can find  $r'$  and  $A'$  if the right sides are known.

Define a new system of coordinates  $x_3 y_3 z_3$  by the next relation, here  $z_3$  axis is directed to OT (see fig. 4) .

	$x_3$	$y_3$	$z_3$	
X	$\cos r' \cos A'$	$-\sin A'$	$\sin r' \cos A'$	(24)
Y	$\cos r' \sin A'$	$\cos A'$	$\sin r' \sin A'$	
Z	$-\sin r'$	0	$\cos r'$	

Now, let the coordinates of F referred to  $(x_3 y_3 z_3)$  system be

$$x_{3F} = OF \sin \theta_4 \cos A_2, \quad y_{3F} = OF \sin \theta_4 \sin A_2, \quad z_{3F} = OF \cos \theta_4 \quad (25)$$

(see fig. 5)

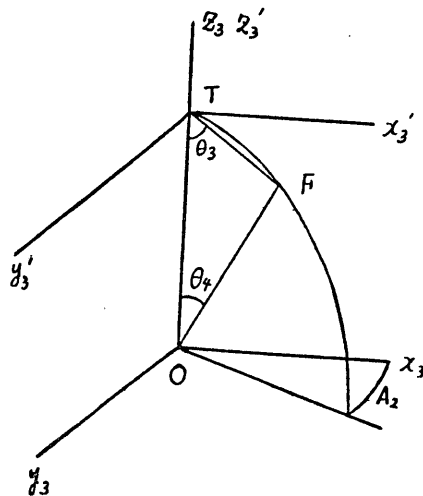


Fig. 5 Relation of  $x_3' y_3' z_3'$  system to  $x_3 y_3 z_3$  system and definition of the angles  $A_2, \theta_3, \theta_4$ .

Further, let  $x_3' y_3' z_3'$  be the parallel translation of  $x_3 y_3 z_3$  system by the translation of O to T (see fig. 6), and the coordinates of F referred to this new system be

$$x_{3'F} = TF \sin \theta_3 \cos A_2, \quad y_{3'F} = TF \sin \theta_3 \sin A_2, \quad z_{3'F} = TF \cos \theta_3 \quad (26)$$

So, if  $\theta_3, A_2$  are given and TF are calculated, the position of F are determined, and in the same way as in T, its height OF can be obtained. From this and

$$\sin \theta_4 = \frac{FT}{OF} \sin \theta_3 \quad (27)$$

we can get  $\theta_4$  and so  $x_{3F} y_{3F} z_{3F}$ , i. e. the coordinate of F referring to  $(x_3 y_3$



$z_3$ ) system, can be determined.

However the coordinates of T referred to  $x_3 y_3 z_3$  system are

$$x_{3T} = y_{3T} = 0, \quad z_{3T} = OT \quad (28)$$

By giving these values we are get the coordinates of F, T referred to X Y Z system  $X_F, Y_F, Z_F, X_T, Y_T, Z_T$  and so

$$X_{F'} = X_F, \quad Y_{F'} = Y_F, \quad Z_{F'} = Z_F - a_0, \quad X_T' = X_T, \quad Y_T' = Y_T, \quad Z_T' = Z_T - a_0 \quad (29)$$

can be determined.

Therefore, the coordinates of T referred to  $X_1, Y_1, Z_1$  system can be obtained as follows :

$$\begin{aligned} r_1 &= (X_T' - X_{F'}) \cosh + (Z_T' - Z_{F'}) \sinh = (X_T - X_F) \cosh + (Z_T - Z_F) \sinh \\ \delta_1 &= Y_T' - Y_{F'} = Y_T - Y_F \\ \kappa_1 &= (Z_T' - Z_{F'}) \cosh - (X_T' - X_{F'}) \sinh = (Z_T - Z_F) \cosh - (X_T - X_F) \sinh \end{aligned} \quad (30)$$

Here,  $h$  is the Sun's altitude and the relation between  $(X, Y, Z_1)$  system and  $(X' Y' Z')$  system is

	$X'$	$Y'$	$Z'$	
$X_1$	$\cosh$	$0$	$\sinh$	
$Y_1$	$0$	$1$	$0$	
$Z_1$	$-\sinh$	$0$	$\cosh$	

(31)

The direction cosines of FT referred to  $X_1, Y_1, Z_1$  axis are

$$\cos \omega_1'' = \frac{r_1}{F T}, \quad \cos \omega_1' = \frac{\delta_1}{F T}, \quad \cos \omega_1 = \frac{\kappa_1}{F T} \quad (32)$$

in which  $\omega_1''$  is the angle between FT and  $X_1$  axis.

$r_1, \delta_1, \kappa_1$  are linear functions of  $x_{3T} - x_{3F}, y_{3T} - y_{3F}, z_{3T} - z_{3F}$  by (17), (24) and (30), and they are expressed as follows by (25)

$$\begin{aligned} x_{3T} - x_{3F} &= x_{3'T} - x_{3'F} = -x_{3'F} = -TF \sin \theta_3 \cos A_2 \\ y_{3T} - y_{3F} &= y_{3'T} - y_{3'F} = -y_{3'F} = -TF \sin \theta_3 \sin A_2 \\ z_{3T} - z_{3F} &= OT - OF \cos \theta_4 = TF \cos \theta_3 \end{aligned} \quad (33)$$

and by (24)

$$\begin{aligned} X_T - X_F &= \cos r' \cos A' (x_{3T} - x_{3F}) - \sin A' (y_{3T} - y_{3F}) + \sin r' \cos A' (z_{3T} - z_{3F}) \\ Y_T - Y_F &= \cos r' \sin A' (x_{3T} - x_{3F}) + \cos A' (y_{3T} - y_{3F}) + \sin r' \sin A' (z_{3T} - z_{3F}) \\ Z_T - Z_F &= -\sin r' (x_{3T} - x_{3F}) + \cos r' (z_{3T} - z_{3F}) \end{aligned} \quad (34)$$

Therefore,  $\cos \omega_1, \cos \omega_1', \cos \omega_1''$ , can be evaluated by sine and cosine of  $r' A' \theta_3 A_2$ . The coordinates of E referred to  $X' Y' Z'$  system are

$$\begin{aligned} X_{1E} &= (X_{E'} - X_{F'}) \cosh + (Z_{E'} - Z_{F'}) \sinh \\ Y_{1E} &= Y_{E'} - Y_{F'} \\ Z_{1E} &= (Z_{E'} - Z_{F'}) \cosh - (X_{E'} - X_{F'}) \sinh \end{aligned} \quad (35)$$

Then

$$X_{1E} - \delta_1 = (X_{E'} - X_T') \cosh + (Z_{E'} - Z_T') \sinh = (X_E - X_T) \cosh + (Z_E - Z_T) \sinh$$

$$Y_{1E} - \delta_1 = Y_{E'} - Y_{T'} = Y_E - Y_T \quad (36)$$

$$Z_{1E} - \kappa_1 = (Z_{E'} - Z_{T'}) \cosh - (X_{E'} - X_{T'}) \sinh = (Z_E - Z_T) \cosh - (X_E - X_T) \sinh$$

The direction cosines of  $x_2 y_2 z_2$  axes with respect to  $X Y Z$  system given in (18) shall be written for brevity by  $a_1, a_2, a_3, b_1, b_2, b_3, c_1, c_2, c_3$ , then

$$\begin{aligned} X_E - X_T &= a_1(x_{2E'} - x_{2T'}) + a_2(y_{2E'} - y_{2T'}) + a_3(z_{2E'} - z_{2T'}) \\ Y_E - Y_T &= b_1(x_{2E'} - x_{2T'}) + b_2(y_{2E'} - y_{2T'}) + b_3(z_{2E'} - z_{2T'}) \\ Z_E - Z_T &= c_1(x_{2E'} - x_{2T'}) + c_2(y_{2E'} - y_{2T'}) + c_3(z_{2E'} - z_{2T'}) \end{aligned} \quad (37)$$

here  $x_{2E'} y_{2E'} z_{2E'}$  being zero by the definition of  $(x_2' y_2' z_2')$  system. Eventually, we get by (18) and (21)

$$\begin{aligned} \frac{X_E - X_T}{E T} &= -\cos r \cos A \sin \theta, \cos A, + \sin A \sin \theta, \sin A, + \sin r \cos A \cos \theta, \\ \frac{Y_E - Y_T}{E T} &= -\cos r \sin A \sin \theta, \cos A, - \cos A \sin \theta, \sin A, + \sin r \sin A \cos \theta, \\ \frac{Z_E - Z_T}{E T} &= \sin r \sin \theta, \cos A, \quad + \cos r \cos \theta, \end{aligned} \quad (38)$$

From this, the direction cosines of TE line referred to  $x, y, z$  system i. e.

$$\frac{X_{1E} - r_1}{T E}, \frac{Y_{1E} - \delta_1}{T E}, \frac{Z_{1E} - \kappa_1}{T E}$$

can be evaluated from (36) by sine and cosine of  $r, A, \theta, A$ . For abridgement we will write them by  $l, m, n$ .

$$l = \frac{X_{1E} - r_1}{T E}, \quad m = \frac{Y_{1E} - \delta_1}{T E}, \quad n = \frac{Z_{1E} - \kappa_1}{T E} \quad (39)$$

Further, let the direction cosines of  $Z_2$  and  $Z_2'$  axes referred to  $X, Y, Z$  system be  $l_2, m_2, n_2$  and  $l_2', m_2', n_2'$ , so we have from (32) and (4)

$$\left. \begin{aligned} \frac{r_1 \kappa_1}{R_1^2} &= \cos \omega, \cos \omega, " = l_2 \sin \omega, = L_1 \\ \frac{\delta_1 \kappa_1}{R_1^2} &= \cos \omega_1' \cos \omega, = m_2 \sin \omega, = M_1 \\ -\frac{r_1^2 + \delta_1^2}{R_1^2} &= -(\cos^2 \omega_1' + \cos^2 \omega, ") = -1 + \cos^2 \omega, = n_2 \sin \omega, = N_1 \end{aligned} \right\} \quad (40)$$

These three quantities shall be notated by  $L_1, M_1$  and  $N_1$ .

$$\cos \omega_2 = l l_2 + m m_2 + n n_2 \quad (41)$$

And, from (32) and (6) we get

$$\begin{aligned} \frac{r_1 \delta_1}{R_1^2} &= \cos \omega, " \cos \omega_1' = l_2 \sin \omega_1', \equiv L_2, \\ -\frac{r_1^2 + \kappa_1^2}{R_1^2} &= -(\cos^2 \omega, + \cos^2 \omega_1, ") - 1 + \cos^2 \omega_1' = m_2' \sin \omega_1', \equiv M_2, \\ \frac{\delta_1 \kappa_1}{R_1^2} &= \cos \omega_1' \cos \omega, = n_2' \sin \omega_1', \equiv N_2, \end{aligned} \quad (42)$$

Three quantities in (42) shall be notated by  $L_2, M_2$  and  $N_2$ .

As  $\omega_2'$  is the angle between TE and  $Z_2'$

$$\cos \omega_2' = l_2' + m_2' + n_2' \quad (43)$$

From (40) and (41)

$$\cos \omega_2 \sin \omega_1 = l_1 + m_1 + n_1 \quad (44)$$

$$\sin^2 \omega_1 \sin^2 \omega_2 = \sin^2 \omega_1 - (l_1 + m_1 + n_1)^2 = \vartheta \quad (45)$$

The expression of (45) shall be for brevity written by  $\vartheta$ .

In the same way

$$\sin^2 \omega_1' \sin^2 \omega_2' = \sin^2 \omega_1' - (l_2 + m_2 + n_2)^2 = \vartheta' \quad (46)$$

The direction cosines of  $X_2$  axis referred to  $X_1, Y_1, Z_1$  system are

$$-\frac{\tau_1}{R_1}, \quad -\frac{\delta_1}{R_1}, \quad -\frac{\kappa_1}{R_1} \quad (47)$$

and that of  $Z_2$  axis  $l_2, m_2, n_2$ , then that of  $Y_2$  can be determined. Let them be  $\lambda, \mu, \nu$ , then we get

	$X_1$	$Y_1$	$Z_1$	
$X_2$	$-\frac{\tau_1}{R_1}$	$\lambda$	$l_2$	(48)
$Y_2$	$-\frac{\delta_1}{R_1}$	$\mu$	$m_2$	
$Z_2$	$-\frac{\kappa_1}{R_1}$	$\nu$	$n_2$	

The result of computation is as follows

$$\lambda = \pm \frac{\delta_1}{\sin \omega_1} \cdot \frac{1}{R_1}$$

$$\mu = \mp \frac{\tau_1}{\sin \omega_1} \cdot \frac{1}{R_1} \quad (49)$$

$$\nu = 0$$

We will explain this result.

It is clear that

$$\lambda \tau_1 + \mu \delta_1 + \nu \kappa_1 = 0 \quad (50)$$

$$\lambda l_2 + \mu m_2 + \nu n_2 = 0 \quad (51)$$

From (50)  $\times l_2 - (51) \times \tau_1$ , we get

$$\mu (\delta_1 l_2 - \tau_1 m_2) + \nu (\kappa_1 l_2 - n_2 \tau_1) = 0 \quad (52)$$

From (50)  $\times m_2 - (51) \times \tau_1$ , we get

$$\lambda (\tau_1 m_2 - \delta_1 l_2) + \nu (\kappa_1 m_2 - \delta_1 n_2) = 0 \quad (53)$$

From (50)  $\times n_2 - (51) \times \kappa_1$ , we get

$$\lambda (\tau_1 n_2 - \kappa_1 l_2) + \mu (\delta_1 n_2 - \kappa_1 m_2) = 0 \quad (54)$$

From (52), (53) and (54), we get

$$\frac{\lambda}{\begin{vmatrix} \delta_1 & \kappa_1 \\ m_2 & n_2 \end{vmatrix}} = \frac{\mu}{\begin{vmatrix} \kappa_1 & \tau_1 \\ n_2 & l_2 \end{vmatrix}} = \frac{\nu}{\begin{vmatrix} \tau_1 & \delta_1 \\ l_2 & m_2 \end{vmatrix}} = -\frac{k}{R_1} \quad (55)$$

Let us define  $k$  by (55). From (40) we get

$$\delta_1 n_2 - \kappa_1 m_2 = -\{\delta_1(\gamma_1^2 + \delta_1^2) + \kappa_1^2 \delta_1\} \frac{1}{R_1^2 \sin \omega_1} = -\frac{\delta_1}{\sin \omega_1} \quad (56)$$

$$\kappa_1 l_2 - n_2 r_1 = \frac{1}{R_1^2 \sin \omega_1} \{\gamma_1 \kappa_1^2 + r_1(\gamma_1^2 + \delta_1^2)\} = \frac{r_1}{\sin \omega_1} \quad (57)$$

$$r_1 m_2 - l_2 \delta_1 = \frac{1}{R_1^2 \sin \omega_1} \{\gamma_1 \delta_1 \kappa_1 - r_1 \kappa_1 \delta_1\} = 0 \quad (58)$$

Substitute (55), (56), (57), (58) in

$$\lambda^2 + \mu^2 + \nu^2 = 1 \quad (59)$$

Then we get

$$\frac{\kappa^2}{R_1^2 \sin^2 \omega_1} (\gamma_1^2 + \delta_1^2) = 1 \quad (60)$$

When we substitute the third expression of (40) in (60), we get

$$\kappa^2 = 1$$

$$\therefore \kappa = \pm 1 \quad (61)$$

Hence we have (49) from (55), (56), (57) and (58).

We may use either the upper or the lower sign in (61) as in explained in later.

Let the direction cosines of  $Y_2'$  axis and  $Z_2'$  axis referring to  $(X, Y, Z_1)$  system be  $\lambda' \mu' \nu'$ ,  $l_2' m_2' n_2'$ . We have then (62).

	$X'_2$	$Y'_2$	$Z'_2$	
$X_1$	$-\frac{r_1}{R_1}$	$\lambda'$	$l_2'$	(62)
$Y_1$	$-\frac{\delta_1}{R_1}$	$\mu'$	$m_2'$	
$Z_1$	$-\frac{\kappa_1}{R_1}$	$\nu'$	$n_2'$	

Therefore, in the same way as in the preceding discussion

$$\frac{\lambda'}{\begin{vmatrix} \delta_1 & \kappa_1 \\ m_2' & n_2' \end{vmatrix}} = \frac{\mu'}{\begin{vmatrix} \kappa_1 & r_1 \\ n_2' & l_2' \end{vmatrix}} = \frac{\nu'}{\begin{vmatrix} r_1 & \delta_1 \\ l_2' & m_2' \end{vmatrix}} = -\frac{k'}{R_1} \quad (63)$$

And from (42) we get

$$\delta_1 n_2' - \kappa_1 m_2' = \frac{1}{R_1^2 \sin \omega_1'} (\delta_1^2 \kappa_1 + \kappa_1 (\gamma_1^2 + \kappa_1^2)) = \frac{\kappa_1}{\sin \omega_1'} \quad (64)$$

$$\kappa_1 l_2' - r_1 n_2' = \frac{1}{R_1^2 \sin \omega_1'} (\kappa_1 r_1 \delta_1 - r_1 \delta_1 \kappa_1) = 0 \quad (65)$$

$$r_1 m_2' - \delta_1 l_2' = \frac{1}{R_1^2 \sin \omega_1'} (-r_1 (\gamma_1^2 + \kappa_1^2) - r_1 \delta_1^2) = -\frac{r_1}{\sin \omega_1'} \quad (66)$$

Substitute (63)~(66) in (67), then we have (68)

$$\begin{aligned}\lambda'^2 + \mu'^2 + \nu'^2 &= 1 \\ \frac{k'^2}{R_1^2} \cdot \frac{1}{\sin^2 \omega_1'} (\tau_1^2 + \kappa_1^2) &= 1 \\ k'^2 &= 1 \\ k' &= \pm 1\end{aligned}\quad (68)$$

Eventually

$$\begin{aligned}\lambda' &= \mp \frac{1}{R_1} \cdot \frac{\kappa_1}{\sin \omega_1'} \\ \mu' &= 0 \\ \nu' &= \pm \frac{1}{R_1} \cdot \frac{\tau_1}{\sin \omega_1'}\end{aligned}\quad (69)$$

Here we can reasonably use each of upper and lower signs.

From (48) we get

$$\begin{aligned}X_1 &= r_1 - \frac{\tau_1}{R_1} X_2 + \lambda Y_2 + l_2 Z_2 \\ Y_1 &= \delta_1 - \frac{\delta_1}{R_1} X_2 + \mu Y_2 + m_2 Z_2 \\ Z_1 &= \kappa_1 - \frac{\kappa_1}{R_1} X_2 + \nu Y_2 + n_2 Z_2\end{aligned}\quad (70)$$

And from (62)

$$\begin{aligned}X_1 &= r_1 - \frac{\tau_1}{R_1} X_2' + \lambda' Y_2' + l_2' Z_2' \\ Y_1 &= \delta_1 - \frac{\delta_1}{R_1} X_2' + \mu' Y_2' + m_2' Z_2' \\ Z_1 &= \kappa_1 - \frac{\kappa_1}{R_1} X_2' + \nu' Y_2' + n_2' Z_2'\end{aligned}\quad (71)$$

Let the directioncosine of  $Z_3$  axis referring to  $(X_2 Y_2 Z_2)$  system be notated by  $(l_3 m_3 n_3)$ , then we have

$$l_3 = \frac{\tau_2 \kappa_2}{R_2^2 \sin \omega_2}, \quad m_3 = \frac{\delta_2 \kappa_2}{R_2^2 \sin \omega_2}, \quad n_3 = -\frac{\tau_2^2 + \delta_2^2}{R_2^2 \sin \omega_2} \quad (72)$$

Here  $(\tau_2 \delta_2 \kappa_2)$  is the coordinate of E referring to  $(X_2 Y_2 Z_2)$  system and can be obtained as follows :

$$\begin{aligned}\tau_2 &= -(X_{1,E} - r_1) \frac{\tau_1}{R_1} - (Y_{1,E} - \delta_1) \frac{\delta_1}{R_1} - (Z_{1,E} - \kappa_1) \frac{\kappa_1}{R_1} \\ \delta_2 &= (X_{1,E} - r_1) \lambda + (Y_{1,E} - \delta_1) \mu + (Z_{1,E} - \kappa_1) \nu \\ \kappa_2 &= (X_{1,E} - r_1) l_2 + (Y_{1,E} - \delta_1) m_2 + (Z_{1,E} - \kappa_1) n_2\end{aligned}$$

Hence we can calculate the above mentioned directioncosines  $(l_3 m_3 n_3)$ .

Let the direction cosine of  $Z_3$  axis referring to  $(X_1 Y_1 Z_1)$  system be  $(\xi \eta \zeta)$ , then we have (74) from (48)

$$\begin{aligned}\xi &= -\frac{r_1}{R_1}l_3 + \lambda m_3 + l_2 n_3 \\ \eta &= -\frac{\delta_1}{R_1}l_3 + \mu m_3 + m_2 n_3 \\ \zeta &= -\frac{\kappa_1}{R_1}l_3 + \nu m_3 + n_2 n_3\end{aligned}\quad (74)$$

However, from (73), (49) and (40)

$$\begin{aligned}\frac{r_2}{R_2} &= -\frac{R_1}{\kappa_1} \left( 1 \frac{r_1 \kappa_1}{R_1^2} + m \frac{\delta_1 \kappa_1}{R_1^2} + n \frac{\kappa_1^2}{R_1^2} \right) = -\frac{R_1}{\kappa_1} (lL_1 + mM_1 + n(N_1 + 1)) \\ &= -\frac{1}{\cos \omega_1} (\sum lL_1 + n)\end{aligned}\quad (75)$$

$$\frac{\delta_2}{R_2} = 1\lambda + m\mu \quad (76)$$

$$\frac{\kappa_2}{R_2} = l_2 + mm_2 + nn_2 = \frac{\sum lL_1}{\sin \omega_1} \quad (77)$$

Then when we define  $L_3$ ,  $M_3$ ,  $N_3$  by (78)

$$\begin{aligned}L_3 &= \frac{r_2 \kappa_2}{R_2^2} \\ M_3 &= \frac{\delta_2 \kappa_2}{R_2^2} \\ N_3 &= -\frac{r_2^2 + \kappa_2^2}{R_2^2}\end{aligned}\quad (78)$$

we have (79) from (72), (75), (76), (77)

$$\begin{aligned}L_3 &= l_3 \sin \omega_2 = -\frac{1}{\sin \omega_1} \cdot \frac{(\sum lL_1)(\sum lL_2 + n)}{\cos \omega_1} \\ M_3 &= m_3 \sin \omega_2 = (1\lambda + m\mu) \frac{\sum lL_1}{\sin \omega_1} \\ N_3 &= n_3 \sin \omega_2 = -\left(1 - \frac{\kappa_2^2}{R_2^2}\right) = -\sin^2 \omega_2 = -\frac{\vartheta}{\sin \omega_1}\end{aligned}\quad (79)$$

Substitute (79) in (74) we get

$$\begin{aligned}\xi &= \frac{1}{\sin \omega_2} \left( -\frac{r_1}{R_1} L_3 + \lambda M_3 + l_2 N_3 \right) \\ \eta &= \frac{1}{\sin \omega_2} \left( -\frac{\delta_1}{R_1} L_3 + \mu M_3 + m_2 N_3 \right) \\ \zeta &= \frac{1}{\sin \omega_2} \left( -\frac{\kappa_1}{R_1} L_3 + \nu M_3 + n_2 N_3 \right)\end{aligned}\quad (80)$$

By further deformation

$$\xi = \frac{1}{\sin \omega_1 \sin \omega_2} \left[ \left\{ \frac{r_1}{R_1} \frac{\sum lL_1 + n}{\cos \omega_1} + \lambda (1\lambda + m\mu) \right\} (\sum lL_1) - \frac{L_1}{\sin^2 \omega_1} \vartheta \right]$$

$$\begin{aligned}\eta &= \frac{1}{\sin \omega_1 \sin \omega_2} \left[ \left\{ \frac{\delta_1}{R_1} \cdot \frac{\sum l_1 + n}{\cos \omega_1} + \mu(1 + m\mu) \right\} (\sum l_1) - \frac{M_1}{\sin^2 \omega_1} \vartheta \right] \\ \zeta &= \frac{1}{\sin \omega_1 \sin \omega_2} \left[ \frac{\kappa_1}{R_1} \cdot \frac{\sum l_1 + n}{\cos \omega_1} \cdot (\sum l_1) - \frac{N_1}{\sin^2 \omega_1} \vartheta \right] \\ &= \frac{1}{\sin \omega_1 \sin \omega_2} \left[ (\sum l_1)(\sum l_1 + n) + \vartheta \right]\end{aligned}\quad (81)$$

Represent the large bracket in (81) by  $\xi, \eta, \zeta$ , then we have

$$\begin{aligned}\xi_1 &= \xi \sin \omega_1 \sin \omega_2 \\ \eta_1 &= \eta \sin \omega_1 \sin \omega_2 \\ \zeta_1 &= \zeta \sin \omega_1 \sin \omega_2\end{aligned}\quad (82)$$

As  $(\xi \eta \zeta)$  is the directioncosine of  $Z_3$  axis referring to  $(X, Y, Z_1)$  system, the value referring to  $(X'Y'Z')$  system reduces to (83) by use of (31)

$$\xi \cosh - \zeta \sinh, \eta, \xi \sinh + \zeta \cosh \quad (83)$$

The directioncosine of  $EO'$  referring to  $(X'Y'Z')$  system is

$$-\cos \theta \cos A, -\cos \theta \sin A, -\sin \theta \quad (84)$$

Thus, as  $\omega_3$  is the angle between  $EO'$  and  $Z_3$  axis, it becomes

$$\begin{aligned}\cos \omega_3 &= -\cos \theta \cos A (\xi \cosh - \zeta \sinh) - \cos \theta \sin A \cdot \eta \\ &\quad - \sin \theta (\xi \sinh + \zeta \cosh)\end{aligned}\quad (85)$$

$$\begin{aligned}\therefore -\sin \omega_1 \sin \omega_2 \cos \omega_3 &= \xi_1 (\sin \theta \sinh + \cos \theta \cosh \cos A) + \eta_1 \cos \theta \sin A \\ &\quad + \zeta_1 (\sin \theta \cosh - \cos \theta \sinh \cos A) = H\end{aligned}\quad (86)$$

If we denote  $H$  as the right hand side of (86), then

$$(1 - \sin^2 \omega_3) \sin^2 \omega_1 \sin^2 \omega_2 = H^2 \quad (87)$$

and when we put

$$D = \sin^2 \omega_1 \sin^2 \omega_2 \sin^2 \omega_3 \quad (88)$$

then we get from (45)  $D = \vartheta - H^2$

When we denote the directioncosine of  $Z_3'$  axis referring to  $(X_2'Y_2'Z_2')$  system as  $(l_3' m_3' n_3')$ , then

$$l_3' = \frac{r_2' \kappa_2'}{R_2' \sin \omega_2'}, \quad m_3' = \frac{\delta_2' \kappa_2'}{R_2' \sin \omega_2'}, \quad n_3' = -\frac{r_2'^2 + \delta_2'^2}{R_2' \sin \omega_2'} \quad (89)$$

Here,  $(r_2' \delta_2' \kappa_2')$  is the coordinate of  $E$  referring to  $(X_2'Y_2'Z_2')$  system and is given as follows by (62)

$$\begin{aligned}r_2' &= -(X_{1E} - r_1) \frac{r_1}{R_1} - (Y_{1E} - \delta_1) \frac{\delta_1}{R_1} - (Z_{1E} - x_1) \frac{\kappa_1}{R_1} \\ \delta_2' &= (X_{1E} - r_1) \lambda' + (Y_{1E} - \delta_1) \mu' + (Z_{1E} - x_1) \nu' \\ \kappa_2' &= (X_{1E} - r_1) I_2' + (Y_{1E} - \delta_1) m_2 + (Z_{1E} - x_1) n_2'\end{aligned}\quad (90)$$

We can then compute the directioncosine  $(l_3' m_3' n_3')$  above mentioned. Then the directioncosine of  $Z_3'$  axis referring to  $(X_1 Y_1 Z_1)$  system can also be computed from (62) as follows

$$\xi' = -\frac{r_1}{R_1} l_3' + \lambda' m_3' + l_2' n_3'$$

$$\eta' = -\frac{\delta_1}{R_1} l_3' + \mu' m_3' + m_2' n_3' \quad (91)$$

$$\zeta' = -\frac{\kappa_1}{R_1} l_3' + \nu' m_3' + n_2' n_3'$$

Let us denote the value as  $(\xi' \eta' \zeta')$ . From (39), (42), (43), (69) and (90) we get

$$\frac{r_2'}{R_2} = \frac{r_2}{R_2} = -\frac{1}{\cos \omega_1} (\sum l L_1 + n) \quad (92)$$

$$\frac{\delta_2'}{R_2} = l \lambda' + m \mu' + n \nu' = l \lambda' + n \nu' \quad (93)$$

$$\frac{\kappa_2'}{R_2} = l l_2' + m m_2' + n n_2' = \frac{\sum l L_2}{\sin \omega_1'} = \cos \omega_2' \quad (94)$$

Then, when we put

$$L_3' = \frac{r_2' \kappa_2'}{R_2^2}$$

$$M_3' = \frac{\delta_2' \kappa_2'}{R_2^2}, \quad N_3' = -\frac{r_2' + \delta_2'}{R_2^2} \quad (95)$$

we get from (89), (92), (93), (94), and (46)

$$\left. \begin{aligned} L_3' &= l_3' \sin \omega_2' = -\frac{\sum l L_1 + n}{\cos \omega_1} \cdot \frac{\sum l L_2}{\sin \omega_1'} \\ M_3' &= m_3' \sin \omega_2' = (l \lambda' + n \nu') \cdot \frac{\sum l L_2}{\sin \omega_1'} \\ N_3' &= n_3' \sin \omega_2' = -\left(1 - \frac{\kappa_2^2}{R_2^2}\right) = -\sin^2 \omega_2' = -\frac{\vartheta'}{\sin^2 \omega_1'} \end{aligned} \right\} \quad (96)$$

Substitute (96) in (91), then we get

$$\xi' = \left(-\frac{r_1}{R_1} L_3' + \lambda M_3' + l_2 N_3'\right) / \sin \omega_2'$$

$$\eta' = \left(-\frac{\delta_1}{R_1} L_3' + \mu' M_3' + m_2' N_3'\right) / \sin \omega_2' \quad (97)$$

$$\zeta' = \left(-\frac{\kappa_1}{R_1} L_3' + \nu' M_3' + n_2' N_3'\right) / \sin \omega_2'$$

By further transformation by (42) and (96)

$$\xi' = \left[ \left\{ \frac{r_1}{R_1} \cdot \frac{(\sum l L_1 + n)}{\cos \omega_1} + \lambda' (l \lambda' + n \nu') \right\} (\sum l L_2) - \frac{L_2}{\sin^2 \omega_1'} \zeta' \right] / \sin \omega_1' \sin \omega_2'$$

$$\eta' = \left[ \left\{ \frac{\delta_1}{R_1} \cdot \frac{(\sum l L_1 + n)}{\cos \omega_1} \cdot (\sum l L_2) - \frac{M_2}{\sin^2 \omega_1'} \zeta' \right\} / \sin \omega_1' \sin \omega_2' \right] \quad (98)$$

$$\zeta' = \left[ \left\{ (\sum l L_1 + n) + \nu' (l \lambda' + n \nu') \right\} (\sum l L_2) - \frac{N_2}{\sin^2 \omega_1'} \vartheta' \right] / \sin \omega_1' \sin \omega_2'$$



When we denote the large bracket in (98) by  $\xi_1', \eta_1', \zeta_1'$ , then we have

$$\begin{aligned}\xi_1' &= \xi' \sin \omega_1' \sin \omega_2' \\ \eta_1' &= \eta' \sin \omega_1' \sin \omega_2' \\ \zeta_1' &= \zeta' \sin \omega_1' \sin \omega_2'\end{aligned}\quad (99)$$

As  $(\xi' \eta' \zeta')$  is the direction cosine of  $Z'_3$  axis referring to  $(X, Y, Z_1)$  system, the value referring to  $(X'Y'Z)$  system is given by (100) from (31).

$$\xi' \cos h - \zeta' \sin h, \quad \eta', \quad \xi' \sin h + \zeta' \cos h \quad (100)$$

As  $\omega_3'$  is the angle between  $EO'$  and  $Z_3'$  axis, then we have from (100) and (84)

$$\begin{aligned}\cos \omega_3' &= -\cos \theta \cos A (\xi' \cos h - \zeta' \sin h) - \cos \theta \sin A \cdot \eta' \\ &\quad - \sin \theta (\xi' \sin h + \zeta' \cos h)\end{aligned}\quad (101)$$

$$\begin{aligned}\therefore -\sin \omega_1' \sin \omega_2' \cos \omega_3' &= (\sin \theta \sin h + \cos \theta \cos h \cos A) \xi_1' \\ &\quad + \cos \theta \sin A \cdot \eta_1' \\ &\quad + (\sin \theta \cos h - \cos \theta \sin h \cos A) \zeta_1' \\ &= H'\end{aligned}\quad (102)$$

We will denote the right hand side of (102) by  $H'$ , then

$$(1 - \sin^2 \omega_3') \sin^2 \omega_1' \sin^2 \omega_2' = H'^2 \quad (103)$$

and when we put

$$D' = \sin^2 \omega_1' \sin^2 \omega_2' \sin^2 \omega_3'$$

then we get from (46)

$$D' = \vartheta' - H'^2 \quad (104)$$

From the above detailed explanation we can evaluate the bracket of (14).

We must now proceed to the discussion of the sign of  $(\lambda \mu \nu)$  in (49). If we adopt the lower sign of (49), it is clear from (73) that the sign of  $\delta_2$  will change, therefore the sign of  $m_3$  will also change from (72), the value of  $(\xi \eta \zeta)$  will be invariable from (72), so that  $(\xi_1, \eta_1, \zeta_1)$  will also be invariable from (81), then  $\cos \omega_3'$  is also so from (85). Moreover  $H$  and  $D$  are invariable from (86) and (88).

In the same way as in this discussion,  $D'$  is invariable either when we adopt the upper sign or the lower sign of (69).

4. Particular case. We can get  $\theta_2$  by giving the right hand side of (20), then  $r'$  and  $A'$  will be gained from (23) by giving  $A, A_1, r$ . By giving  $\theta_3, A_2$  we can get  $r_1/R_1, \delta_1/R_1, \kappa_1/R_1$  from (33), (34) and (30).

We can get  $l, m, n$ , from (38) and (36), and  $L, M, N, l_2 m_2 n_2$  from (40),  $\omega_2$  from (41),  $L_2 M_2 N_2 l_2' m_2' n_2'$  from (42),  $\omega_2'$  from (45),  $\vartheta$  from (45)  $\vartheta'$  from (46),  $(\lambda, \mu, \nu)$  from (49),  $(\xi_1, \eta_1, \zeta_1)$  from (81),  $H$  from (86) and eventually  $D$  from (88).

In the same way, we can get  $(\lambda' \mu' \nu')$  from (69),  $(\xi_1', \eta_1', \zeta_1')$  from (98),  $H'$  from (102) and at last  $D'$  from (104).

To simplify the discussion, let us put

$$\frac{X_T - X_F}{TF} = a, \quad \frac{Y_T - Y_F}{TF} = b, \quad \frac{Z_T - Z_F}{TF} = c \quad (105)$$

then we have from (33) and (34)

a) The case  $A_2 = 0$

$$\begin{aligned} a &= \cos A' \sin(\tau' - \theta_3) \\ b &= \sin A' \sin(\tau' - \theta_3) \\ c &= \cos(\tau' - \theta_3) \end{aligned} \quad (106)$$

b) The case  $A_2 = \frac{\pi}{2}$

$$\begin{aligned} a &= \sin A' \sin \theta_3 + \sin \tau' \cos A' \cos \theta_3 \\ b &= -\cos A' \sin \theta_3 + \sin \tau' \sin A' \cos \theta_3 \\ c &= \cos \tau' \cos \theta_3 \end{aligned} \quad (107)$$

c) The case  $A_2 = \pi$

$$\begin{aligned} a &= \cos A' \sin(\tau' + \theta_3) \\ b &= \sin A' \sin(\tau' + \theta_3) \\ c &= \cos(\tau' + \theta_3) \end{aligned} \quad (108)$$

d) The case  $A_2 = \frac{3\pi}{2}$

$$\begin{aligned} a &= -\sin A' \sin \theta_3 + \sin \tau' \cos A' \cos \theta_3 \\ b &= \cos A' \sin \theta_3 + \sin \tau' \sin A' \cos \theta_3 \\ c &= \cos \tau' \cos \theta_3 \end{aligned} \quad (109)$$

We have further from (23)

I) The case  $A = 0$

$$\begin{aligned} \sin \tau' \cos A' &= \cos \tau \sin \theta_2 \cos A_1 + \sin \tau \cos \theta_2 \\ \sin \tau' \sin A' &= \sin \theta_2 \sin A_1 \\ \cos \tau &= -\sin \tau \sin \theta_2 \cos A_1 + \cos \tau \cos \theta_2 \end{aligned} \quad (110)$$

1) The case  $A_1 = 0$

$$\begin{aligned} \sin \tau' \cos A' &= \sin(\tau + \theta_2) \\ \sin \tau' \sin A' &= 0 \\ \cos \tau' &= \cos(\tau + \theta_2) \\ \therefore A &= 0, \quad \tau' = \tau + \theta_2 \end{aligned} \quad (111)$$

then we get from (106)~(109)

$$\begin{aligned} \text{a) } a &= \sin(\tau + \theta_2 - \theta_3) \\ b &= 0 \\ c &= \cos(\tau + \theta_2 - \theta_3) \end{aligned} \quad (112)$$

And we get from (30)

$$\left. \begin{aligned} \frac{r_1}{R_1} &= \sin(\tau + \theta_2 - \theta_3 + h) \\ \frac{\delta_1}{R_1} &= 0 \\ \frac{\kappa_1}{R_1} &= \cos(\tau + \theta_2 - \theta_3 + h) \end{aligned} \right\} \quad (113)$$

$$\begin{aligned}
 \text{b) } & \left. \begin{aligned} a &= \sin r' \cos \theta_3 \\ b &= -\sin \theta_3 \\ c &= \cos r' \cos \theta_3 \end{aligned} \right\} \quad (114)
 \end{aligned}$$

$$\left. \begin{aligned} \frac{r_1}{R_1} &= \cos \theta_3 \sin(\gamma + \theta_2 + h) \\ \frac{\delta_1}{R_1} &= -\sin \theta_3 \\ \frac{\kappa_1}{R_1} &= \cos \theta_3 \cos(\gamma + \theta_2 + h) \end{aligned} \right\} \quad (115)$$

$$\begin{aligned}
 \text{c) } & \left. \begin{aligned} a &= \sin(\gamma' + \theta_3) \\ b &= 0 \\ c &= \cos(\gamma' + \theta_3) \end{aligned} \right\} \quad (116)
 \end{aligned}$$

$$\left. \begin{aligned} \frac{r_1}{R_1} &= \sin(\gamma + \theta_2 + \theta_3 + h) \\ \frac{\delta_1}{R_1} &= 0 \\ \frac{\kappa_1}{R_1} &= \cos(\gamma + \theta_2 + \theta_3 + h) \end{aligned} \right\} \quad (117)$$

$$\begin{aligned}
 \text{d) } & \left. \begin{aligned} a &= \sin r' \cos \theta_3 \\ b &= \sin \theta_3 \\ c &= \cos r' \cos \theta_3 \end{aligned} \right\} \quad (118)
 \end{aligned}$$

$$\left. \begin{aligned} \frac{r_1}{R_1} &= \cos \theta_3 \sin(\gamma + \theta_2 + h) \\ \frac{\delta_1}{R_1} &= \sin \theta_3 \\ \frac{\kappa_1}{R_1} &= \cos \theta_3 \cos(\gamma + \theta_2 + h) \end{aligned} \right\} \quad (119)$$

In this paper the values of  $D D'$  in the case I) 1) c) are tabulated in Table 1 for  $r = 0$ , here the value of  $\theta, \theta_1, \theta_2, \theta_3$  and  $h$  are expressed in unit of degree. In this case  $D + D' = 1 + \vartheta \sin^2(\theta - \theta_1)$ . We have the values in the case of I) 1) a) when we wright the column of  $\theta_3$  upside down i. e. substitute  $\pi - \theta_3$  for  $\theta_3$ . At the end the author expresses his sincere thanks to Dr. K. Y. Kondratyev, Rector of Leningrad University, for encouragement and interest to the research.

**References**

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Table 1.

$\theta = 0$

This table expresses the value of the polarization angle in the case of I) 1) c)

$\theta_1$	$\theta_2$ $\theta_3$	h																			
		0					30					60					90				
		0	1	2	3	4	0	1	2	3	4	0	1	2	3	4	0	1	2	3	4
30	0	1.0000	1.0000	1.0002	1.0004	1.0008	1.0468	1.0487	1.0505	1.0521	1.0537	1.1406	1.1404	1.1402	1.1396	1.1388	1.1875	1.1836	1.1795	1.1755	1.1709
	30	1.0156	1.0155	1.0154	1.0152	1.0150	1.0468	1.0449	1.0429	1.0409	1.0388	1.0625	1.0588	1.0550	1.0513	1.0478	1.0468	1.0431	1.0396	1.0362	1.0330
	60	1.0000	1.0000	1.0002	1.0005	1.0009	1.0000	1.0000	1.0003	1.0007	1.0012	1.0000	1.0001	1.0002	1.0005	1.0008	1.0000	1.0000	1.0000	1.0001	1.0002
	75	1.0156	1.0179	1.0202	1.0228	1.0255	1.0156	1.0175	1.0195	1.0215	1.0237	1.0083	1.0091	1.0099	1.0107	1.0114	1.0011	1.0011	1.0010	1.0010	1.0009
	90	1.0625	1.0663	1.0701	1.0740	1.0778	1.0468	1.0487	1.0504	1.0521	1.0537	1.0156	1.0156	1.0154	1.0152	1.0150	1.0000	1.0000	1.0000	1.0004	1.0004
	105	1.1166	1.1195	1.1222	1.1248	1.1273	1.0625	1.0624	1.0622	1.0618	1.0613	1.0083	1.0075	1.0067	1.0059	1.0051	1.0084	1.0099	1.0114	1.0132	1.0151
	120	1.1406	1.1405	1.1401	1.1396	1.1388	1.0468	1.0449	1.0429	1.0409	1.0388	1.0000	1.0000	1.0002	1.0005	1.0010	1.0468	1.0507	1.0547	1.0588	1.0631
	150	1.0625	1.0587	1.0550	1.0513	1.0478	1.0000	1.0000	1.0003	1.0006	1.0012	1.0625	1.0662	1.0701	1.0739	1.0778	1.1875	1.1911	1.1947	1.1979	1.2009
180												1.1404									
60	0	1.0000	1.0000	1.0000	1.0003	1.0007	1.0468	1.0467	1.0464	1.0458	1.0450	1.1406	1.1348	1.1288	1.1227	1.1164	1.1875	1.1764	1.1650	1.1541	1.1434
	30	1.0000	1.0000	1.0000	1.0006	1.0010	1.0000	1.0002	1.0006	1.0016	1.0029	1.0000	1.0001	1.0009	1.0021	1.0036	1.0000	1.0001	1.0007	1.0016	1.0024
	60	1.1406	1.1521	1.1641	1.1765	1.1894	1.1875	1.1989	1.2103	1.2220	1.2354	1.1406	1.1461	1.1514	1.1564	1.1612	1.0468	1.0467	1.0463	1.0459	1.0451
	75	1.3498	1.3663	1.3808	1.3963	1.4116	1.3498	1.3585	1.3668	1.3746	1.3820	1.1875	1.1872	1.1866	1.1854	1.1839	1.0251	1.0226	1.0203	1.0179	1.0155
	90	1.5625	1.5736	1.5840	1.5937	1.6030	1.4218	1.4215	1.4205	1.4188	1.4164	1.1406	1.1348	1.1288	1.1227	1.1164	1.0000	1.0001	1.0006	1.0015	1.0030
	105	1.6528	1.6525	1.6511	1.6490	1.6461	1.3498	1.3408	1.3312	1.3213	1.3111	1.0468	1.0413	1.0360	1.0309	1.0243	1.0468	1.0537	1.0609	1.0685	1.0765
	120	1.5625	1.5509	1.5388	1.5261	1.5128	1.1875	1.1761	1.1650	1.1541	1.1434	1.0000	1.0002	1.0009	1.0020	1.0025	1.1875	1.1988	1.2103	1.2218	1.2334
	150	1.1406	1.1294	1.1188	1.1087	1.0990	1.0000	1.0001	1.0006	1.0014	1.0025	1.1406	1.1461	1.1515	1.1564	1.0450	1.4218	1.0056	1.4206	1.4188	1.4164
180																					
75	0	1.0000	1.0000	1.0000	1.0000	1.0001	1.0156	1.0145	1.0133	1.0119	1.0107	1.0469	1.0418	1.0367	1.0320	1.0273	1.0620	1.0545	1.0473	1.0405	1.0337
	30	1.0156	1.0189	1.0222	1.0264	1.0309	1.0470	1.0543	1.0624	1.0709	1.0800	1.0625	1.0708	1.0796	1.0891	1.0983	1.0469	1.0522	1.0572	1.0626	1.0680
	60	1.3499	1.3692	1.3889	1.4090	1.4292	1.4666	1.4826	1.4985	1.5140	1.5290	1.3499	1.3551	1.3588	1.3624	1.3652	1.1166	1.1134	1.1100	1.1061	1.1021
	75	1.6528	1.6719	1.6906	1.7088	1.7261	1.6530	1.6595	1.6652	1.6731	1.6739	1.3498	1.3444	1.3383	1.3316	1.3244	1.0469	1.0417	1.0367	1.0320	1.0274
	90	1.8704	1.8782	1.8847	1.8901	1.8946	1.6528	1.6454	1.6371	1.6279	1.6178	1.2176	1.2063	1.1951	1.1839	1.1727	1.0000	1.0002	1.0010	1.0024	1.0043
	105	1.8704	1.8619	1.8522	1.8416	1.8299	1.4665	1.4501	1.4333	1.4165	1.3996	1.0625	1.0545	1.0471	1.0402	1.0337	1.0625	1.0709	1.0796	1.0888	1.0984
	120	1.6529	1.6336	1.6137	1.5935	1.5731	1.2176	1.2025	1.1879	1.1739	1.1602	1.0000	1.0002	1.0010	1.0022	1.0041	1.2176	1.2285	1.2395	1.2503	1.2608
	150	1.1166	1.1058	1.0955	1.0861	1.0771	1.0000	1.0000	1.0004	1.0011	1.0019	1.2176	1.2285	1.2395	1.2503	1.2608	1.3499	1.3443	1.3383	1.3316	1.3244
180																					
90	0	1.0000	1.0000	1.0000	1.0000	1.0001	1.0000	1.0001	1.0003	1.0008	1.0014	1.0000	1.0003	1.0009	1.0022	1.0039	1.0000	1.0002	1.0012	1.0028	1.0048
	30	1.0625	1.0704	1.0788	1.0879	1.0978	1.1877	1.2028	1.2188	1.2354	1.2526	1.2500	1.2652	1.2805	1.2960	1.3113	1.1875	1.1949	1.2019	1.2086	1.2150
	60	1.5625	1.5851	1.6078	1.6303	1.6526	1.7500	1.7648	1.7787	1.7917	1.8040	1.5335	1.5621	1.5607	1.5585	1.5552	1.1875	1.1798	1.1718	1.1636	1.1553
	75	1.8705	1.8864	1.9015	1.9155	1.9284	1.8705	1.8700	1.8682	1.8654	1.8615	1.4665	1.4544	1.4416	1.4284	1.4148	1.0625	1.0551	1.0480	1.0413	1.0351
	90	2.0000	1.9995	1.9976	1.9943	1.9904	1.7500	1.7346	1.7184	1.7015	1.6838	1.2500	1.2349	1.2201	1.2055	1.1913	1.0000	1.0003	1.0012	1.0027	1.0049
	105	1.8705	1.8539	1.8363	1.8183	1.7992	1.4665	1.4460	1.4253	1.4050	1.3849	1.0625	1.0541	1.0463	1.0391	1.0325	1.0625	1.0702	1.0782	1.0864	1.0948
	120	1.5625	1.5400	1.5173	1.4947	1.4723	1.1875	1.1726	1.1585	1.1450	1.1321	1.0000	1.0002	1.0009	1.0019	1.0034	1.1875	1.1949	1.2017	1.2086	1.2149
	150	1.0625	1.0552	1.0485	1.0425	1.0370	1.0000	1.0001	1.0003	1.0005	1.0010	1.0335	1.0623	1.0619	1.0611	1.0601	1.1875	1.1798	1.1718	1.1636	1.1553

105	0	1.0000	1.0000	1.0000	1.0001	1.0005	1.0156	1.0189	1.0223	1.0264	1.0309	1.0470	1.0543	1.0624	1.0709	1.0800	1.0620	1.0708	1.0796	1.0891	1.0983
	30	1.1166	1.1281	1.1401	1.1528	1.1665	1.3499	1.3692	1.3889	1.4090	1.4292	1.4666	1.4826	1.4985	1.5140	1.5290	1.3499	1.3547	1.3588	1.3624	1.3652
	60	1.6529	1.6719	1.6906	1.7087	1.7262	1.8704	1.8782	1.8847	1.8901	1.8946	1.6528	1.6454	1.6371	1.6279	1.6178	1.2176	1.2063	1.1951	1.1839	1.1727
	75	1.8704	1.8781	1.8847	1.8903	1.8957	1.8704	1.8619	1.8522	1.8416	1.8299	1.4665	1.4501	1.4333	1.4165	1.3996	1.0625	1.0545	1.0471	1.0402	1.0337
	90	1.8704	1.8619	1.8522	1.8415	1.8301	1.6529	1.6336	1.6137	1.5935	1.5731	1.2176	1.2025	1.1879	1.1739	1.1602	1.0000	1.0002	1.0010	1.0022	1.0041
	105	1.6528	1.6335	1.6136	1.5936	1.5732	1.3498	1.3308	1.3120	1.2938	1.2760	1.0469	1.0401	1.0339	1.0281	1.0233	1.0469	1.0520	1.0573	1.0626	1.0680
	120	1.3499	1.3309	1.3120	1.2938	1.2759	1.1166	1.1058	1.0955	1.0861	1.0771	1.0000	1.0000	1.0004	1.0011	1.0019	1.1166	1.1194	1.1218	1.1239	1.1255
150	1.0156	1.0128	1.0103	1.0083	1.0065	1.0000	1.0000	1.0000	1.0000	1.0001	1.0156	1.0145	1.0133	1.0119	1.0106	1.0469	1.0418	1.0367	1.0320	1.0273	
120	0	1.0000	1.0000	1.0002	1.0006	1.0012	1.0468	1.0528	1.0591	1.0659	1.0733	1.1406	1.1521	1.1641	1.1765	1.1894	1.1375	1.1989	1.2103	1.2220	1.2334
	30	1.1406	1.1521	1.1641	1.1765	1.1894	1.4218	1.4388	1.4558	1.4727	1.4894	1.5625	1.5736	1.5840	1.5937	1.6034	1.4219	1.4215	1.4205	1.4188	1.4164
	60	1.5625	1.5735	1.5840	1.5937	1.6029	1.7500	1.7496	1.7482	1.7458	1.7428	1.5625	1.5509	1.5388	1.5261	1.5128	1.3375	1.1761	1.1650	1.1541	1.1434
	75	1.6528	1.6524	1.6512	1.6491	1.6460	1.6528	1.6404	1.6272	1.6137	1.5994	1.3498	1.3345	1.3189	1.3037	1.2886	1.0468	1.0405	1.0347	1.0293	1.0243
	90	1.5625	1.5509	1.5387	1.5260	1.5129	1.4218	1.4050	1.3879	1.3710	1.3542	1.1406	1.1294	1.1188	1.1087	1.0990	1.0000	1.0001	1.0006	1.0014	1.0025
	105	1.3498	1.3344	1.3189	1.3039	1.2886	1.1875	1.1746	1.1622	1.1503	1.1389	1.0251	1.0211	1.0176	1.0144	1.0117	1.0251	1.0275	1.0298	1.0320	1.0342
	120	1.1406	1.1295	1.1188	1.1086	1.0990	1.0468	1.0414	1.0363	1.0318	1.0277	1.0000	1.0000	1.0002	1.0003	1.0007	1.0468	1.0467	1.0464	1.0458	1.0450
150	1.0000	1.0000	1.0002	1.0003	1.0007	1.0000	1.0000	1.0000	1.0000	1.0000	1.0000	1.0000	1.0002	1.0006	1.0010	1.0000	1.0002	1.0006	1.0012	1.0029	
150	0	1.0000	1.0000	1.0002	1.0006	1.0010	1.0468	1.0507	1.0547	1.0588	1.0631	1.1406	1.1462	1.1519	1.1575	1.1631	1.1875	1.1912	1.1946	1.1979	1.2010
	30	1.0625	1.0662	1.0701	1.0739	1.0778	1.1875	1.1911	1.1946	1.1979	1.2009	1.2500	1.2498	1.2494	1.2486	1.2476	1.1875	1.1836	1.1796	1.1753	1.1709
	60	1.1406	1.1405	1.1402	1.1396	1.1388	1.1875	1.1836	1.1795	1.1753	1.1709	1.1406	1.1350	1.1293	1.1236	1.1180	1.0468	1.0431	1.0396	1.0362	1.0330
	75	1.1166	1.1136	1.1104	1.1071	1.1037	1.1166	1.1114	1.1063	1.1013	1.0962	1.0625	1.0582	1.0540	1.0501	1.0463	1.0083	1.0070	1.0058	1.0048	1.0039
	90	1.0625	1.0587	1.0550	1.0513	1.0478	1.0468	1.0431	1.0396	1.0362	1.0330	1.0156	1.0138	1.0121	1.0106	1.0092	1.0000	1.0000	1.0000	1.0001	1.0002
	105	1.0156	1.0135	1.0115	1.0097	1.0081	1.0083	1.0070	1.0058	1.0048	1.0039	1.0011	1.0008	1.0006	1.0004	1.0003	1.0011	1.0011	1.0010	1.0010	1.0009
	120	1.0000	1.0001	1.0002	1.0005	1.0008	1.0000	1.0000	1.0000	1.0000	1.0001	1.0002	1.0000	1.0000	1.0000	1.0000	1.0000	1.0000	1.0000	1.0000	1.0003
150	1.0156	1.0156	1.0154	1.0152	1.0150	1.0000	1.0000	1.0000	1.0002	1.0004	1.0156	1.0176	1.0197	1.0219	1.0244	1.0468	1.0507	1.0547	1.0588	1.0631	

The Phase function in the Tertiary Scattering (1)

Table 1. A = 0  $\theta = 30$

$\theta_1$	$\theta_2$ $\theta_3$	0					30					60					90				
		0	1	2	3	4	0	1	2	3	4	0	1	2	3	4	0	1	2	3	4
		h																			
30	0																				
	30																				
	60																				
	75																				
	90	1.0000																			
	105																				
	180																				
60	0	1.0000	1.0000	1.0001	1.0001	1.0003	1.0156	1.0156	1.0155	1.0153	1.0150	1.0469	1.0450	1.0430	1.0409	1.0388	1.0625	1.0588	1.0550	1.0514	1.0478
	30	1.0000	1.0000	1.0001	1.0002	1.0004	1.0000	1.0001	1.0002	1.0006	1.0010	1.0000	1.0001	1.0003	1.0007	1.0012	1.0000	1.0001	1.0003	1.0006	1.0008
	60	1.0469	1.0507	1.0547	1.0589	1.0632	1.0625	1.0663	1.0701	1.0740	1.0778	1.0469	1.0487	1.0505	1.0522	1.0538	1.0156	1.0156	1.0155	1.0153	1.0151
	75	1.1166	1.1218	1.1270	1.1321	1.1372	1.1166	1.1195	1.1223	1.1249	1.1274	1.0625	1.0624	1.0622	1.0618	1.0613	1.0084	1.0076	1.0068	1.0060	1.0052
	90	1.1875	1.1912	1.1947	1.1979	1.2010	1.1406	1.1405	1.1402	1.1396	1.1388	1.0469	1.0450	1.0430	1.0409	1.0388	1.0000	1.0001	1.0002	1.0005	1.0010
	105	1.2176	1.2175	1.2171	1.2164	1.2154	1.1166	1.1136	1.1104	1.1071	1.1037	1.0156	1.0138	1.0120	1.0103	1.0088	1.0156	1.0179	1.0203	1.0229	1.0255
	180	1.1875	1.1837	1.1796	1.1754	1.1710	1.0625	1.0587	1.0550	1.0514	1.0478	1.0000	1.0001	1.0003	1.0007	1.0012	1.0625	1.0663	1.0701	1.0740	1.0778
	150	1.0469	1.0432	1.0396	1.0363	1.0330	1.0000	1.0001	1.0002	1.0005	1.0009	1.0469	1.0487	1.0505	1.0522	1.0537	1.1406	1.1405	1.1402	1.1396	1.1388
												1.0469	1.0450	1.0430	1.0409	1.0388	1.0625	1.0588	1.0550	1.0514	1.0478
75	0	1.0000	1.0000	1.0000	1.0001	1.0001	1.0084	1.0078	1.0072	1.0064	1.0058	1.0252	1.0225	1.0197	1.0172	1.0147	1.0333	1.0293	1.0254	1.0218	1.0181
	30	1.0084	1.0102	1.0119	1.0142	1.0166	1.0252	1.0292	1.0335	1.0380	1.0437	1.0335	1.0380	1.0427	1.0478	1.0527	1.0252	1.0280	1.0307	1.0336	1.0365
	60	1.1876	1.1979	1.2085	1.2192	1.2301	1.2501	1.2587	1.2672	1.2755	1.2835	1.1876	1.1904	1.1923	1.1943	1.1958	1.0625	1.0608	1.0590	1.0569	1.0548
	75	1.3499	1.3601	1.3702	1.3799	1.3892	1.3500	1.3535	1.3565	1.3608	1.3612	1.1875	1.1846	1.1813	1.1778	1.1759	1.0252	1.0224	1.0197	1.0172	1.0147
	90	1.4665	1.4706	1.4742	1.4771	1.4795	1.3499	1.3459	1.3415	1.3365	1.3311	1.1167	1.1106	1.1046	1.0986	1.0926	1.0000	1.0002	1.0006	1.0013	1.0024
	105	1.4665	1.4619	1.4567	1.4511	1.4448	1.2530	1.2413	1.2323	1.2233	1.2142	1.0335	1.0293	1.0253	1.0216	1.0181	1.0335	1.0380	1.0427	1.0476	1.0528
	180	1.3499	1.3396	1.3289	1.3181	1.3072	1.1167	1.1086	1.1008	1.0932	1.0859	1.0000	1.0002	1.0006	1.0012	1.0022	1.1167	1.1225	1.1284	1.1342	1.1398
	150	1.0625	1.0567	1.0512	1.0462	1.0414	1.0000	1.0001	1.0003	1.0006	1.0011	1.1167	1.1225	1.1284	1.1342	1.1398	1.1876	1.1846	1.1813	1.1778	1.1759
												1.0252	1.0225	1.0197	1.0172	1.0147	1.0333	1.0293	1.0254	1.0218	1.0181
90	0	1.0000	1.0000	1.0000	1.0000	1.0001	1.0000	1.0001	1.0002	1.0006	1.0011	1.0000	1.0002	1.0007	1.0017	1.0029	1.0000	1.0002	1.0009	1.0021	1.0036
	30	1.0469	1.0528	1.0591	1.0659	1.0734	1.1408	1.1521	1.1641	1.1766	1.1895	1.1875	1.1989	1.2104	1.2220	1.2335	1.1406	1.1462	1.1514	1.1565	1.1613
	60	1.4219	1.4388	1.4559	1.4727	1.4896	1.5625	1.5736	1.5840	1.5938	1.6030	1.4001	1.4216	1.4205	1.4189	1.4164	1.1406	1.1349	1.1289	1.1227	1.1165
	75	1.6529	1.6648	1.6761	1.6866	1.6963	1.6529	1.6525	1.6512	1.6491	1.6461	1.3499	1.3408	1.3312	1.3213	1.3111	1.0469	1.0413	1.0360	1.0310	1.0263
	90	1.7500	1.7496	1.7482	1.7459	1.7428	1.5625	1.5510	1.5388	1.5261	1.5129	1.1875	1.1762	1.1651	1.1541	1.1435	1.0000	1.0002	1.0009	1.0020	1.0037
	105	1.6529	1.6404	1.6272	1.6137	1.5994	1.3499	1.3345	1.3190	1.3038	1.2887	1.0469	1.0406	1.0347	1.0293	1.0244	1.0469	1.0527	1.0587	1.0648	1.0711
	180	1.4219	1.4050	1.3880	1.3710	1.3542	1.1406	1.1295	1.1189	1.1088	1.0991	1.0000	1.0002	1.0007	1.0014	1.0026	1.1406	1.1462	1.1513	1.1565	1.1612
	150	1.0469	1.0414	1.0364	1.0319	1.0278	1.0000	1.0001	1.0002	1.0004	1.0008	1.0251	1.0467	1.0464	1.0458	1.0451	1.1406	1.1349	1.1289	1.1227	1.1165
												1.0000	1.0002	1.0007	1.0017	1.0029	1.0000	1.0002	1.0009	1.0021	1.0036

105	0	1.0000	1.0000	1.0000	1.0001	1.0005	1.0156	1.0189	1.0223	1.0264	1.0309	1.0470	1.0544	1.0624	1.0709	1.0801	1.0620	1.0708	1.0797	1.0891	1.0983
	30	1.1166	1.1281	1.1401	1.1528	1.1665	1.3499	1.3692	1.3889	1.4090	1.4292	1.4667	1.4826	1.4986	1.5141	1.5290	1.3500	1.3547	1.3588	1.3625	1.3653
	60	1.6529	1.6719	1.6906	1.7087	1.7262	1.8704	1.8782	1.8847	1.8901	1.8946	1.6528	1.6454	1.6371	1.6279	1.6178	1.2177	1.2064	1.1952	1.1840	1.1728
	75	1.8704	1.8781	1.8847	1.8903	1.8967	1.8704	1.8619	1.8522	1.8416	1.8299	1.4665	1.4502	1.4354	1.4166	1.3997	1.0625	1.0546	1.0472	1.0402	1.0338
	90	1.8704	1.8619	1.8522	1.8415	1.8301	1.6529	1.6336	1.6137	1.5935	1.5731	1.2177	1.2026	1.1880	1.1739	1.1603	1.0000	1.0003	1.0010	1.0022	1.0041
	105	1.6528	1.6335	1.6136	1.5936	1.5732	1.3498	1.3308	1.3120	1.2938	1.2760	1.0469	1.0401	1.0340	1.0282	1.0233	1.0469	1.0521	1.0574	1.0627	1.0680
	120	1.3499	1.3309	1.3120	1.2938	1.2759	1.1166	1.1058	1.0955	1.0861	1.0771	1.0000	1.0001	1.0005	1.0011	1.0020	1.1166	1.1194	1.1218	1.1239	1.1256
150	1.0156	1.0128	1.0103	1.0083	1.0065	1.0000	1.0000	1.0000	1.0000	1.0001	1.0157	1.0146	1.0135	1.0119	1.0106	1.0470	1.0419	1.0368	1.0321	1.0275	
											1.0470	1.0544	1.0624	1.0709	1.0801	1.0620	1.0708	1.0797	1.0891	1.0983	
120	0	1.0000	1.0001	1.0003	1.0008	1.0016	1.0625	1.0704	1.0788	1.0879	1.0978	1.1875	1.2028	1.2188	1.2354	1.2526	1.2500	1.2652	1.2805	1.2960	1.3113
	30	1.1875	1.2028	1.2189	1.2354	1.2526	1.5625	1.5851	1.6078	1.6303	1.6526	1.7600	1.7648	1.7787	1.7917	1.8040	1.5625	1.5621	1.5607	1.5585	1.5552
	60	1.7500	1.7647	1.7787	1.7917	1.8039	2.0000	1.9995	1.9976	1.9945	1.9904	1.7500	1.7346	1.7184	1.7015	1.6838	1.2500	1.2349	1.2201	1.2055	1.1913
	75	1.8705	1.8699	1.8683	1.8655	1.8614	1.8705	1.8539	1.8363	1.8183	1.7992	1.4665	1.4460	1.4253	1.4050	1.3849	1.0625	1.0541	1.0463	1.0391	1.0325
	90	1.7500	1.7346	1.7183	1.7014	1.6839	1.5625	1.5400	1.5173	1.4947	1.4723	1.1875	1.1726	1.1585	1.1450	1.1321	1.0000	1.0002	1.0009	1.0019	1.0034
	105	1.4665	1.4459	1.4253	1.4052	1.3848	1.2500	1.2329	1.2163	1.2005	1.1853	1.0335	1.0282	1.0235	1.0193	1.0157	1.0335	1.0367	1.0398	1.0427	1.0456
	120	1.1875	1.1727	1.1584	1.1449	1.1320	1.0625	1.0552	1.0485	1.0425	1.0370	1.0000	1.0001	1.0003	1.0005	1.0010	1.0625	1.0623	1.0619	1.0611	1.0601
150	1.0000	1.0001	1.0003	1.0005	1.0010	1.0000	1.0000	1.0000	1.0000	1.0001	1.0000	1.0001	1.0003	1.0008	1.0014	1.0000	1.0003	1.0009	1.0017	1.0039	
											1.1875	1.2028	1.2188	1.2354	1.2526	1.2500	1.2652	1.2805	1.2960	1.3113	
150	0	1.0000	1.0002	1.0007	1.0020	1.0030	1.1406	1.1521	1.1642	1.1766	1.1895	1.4219	1.4388	1.4559	1.4727	1.4895	1.5625	1.5736	1.5840	1.5938	1.6030
	30	1.1875	1.1988	1.2104	1.2219	1.2334	1.5625	1.5735	1.5840	1.5938	1.6029	1.7500	1.7496	1.7482	1.7459	1.7428	1.5625	1.5510	1.5388	1.5261	1.5129
	60	1.4219	1.4215	1.4206	1.4189	1.4164	1.5625	1.5510	1.5387	1.5261	1.5129	1.4219	1.4050	1.3880	1.3710	1.3542	1.1406	1.1295	1.1189	1.1088	1.0991
	75	1.3499	1.3408	1.3312	1.3213	1.3112	1.3499	1.3344	1.3190	1.3039	1.2886	1.1875	1.1747	1.1622	1.1504	1.1390	1.0251	1.0212	1.0176	1.0145	1.0118
	90	1.1875	1.1763	1.1650	1.1541	1.1435	1.1406	1.1295	1.1188	1.1087	1.0990	1.0469	1.0414	1.0364	1.0319	1.0278	1.0000	1.0001	1.0002	1.0004	1.0008
	105	1.0469	1.0405	1.0347	1.0293	1.0243	1.0251	1.0212	1.0176	1.0145	1.0117	1.0034	1.0026	1.0020	1.0014	1.0010	1.0034	1.0034	1.0032	1.0031	1.0029
	120	1.0000	1.0003	1.0007	1.0015	1.0026	1.0000	1.0001	1.0002	1.0004	1.0008	1.0000	1.0000	1.0000	1.0000	1.0001	1.0000	1.0001	1.0002	1.0006	1.0011
150	1.0469	1.0468	1.0464	1.0458	1.0451	1.0000	1.0001	1.0002	1.0006	1.0012	1.0469	1.0528	1.0591	1.0659	1.0734	1.1406	1.1521	1.1641	1.1766	1.1895	
											1.4219	1.4388	1.4559	1.4727	1.4895	1.5625	1.5736	1.5840	1.5938	1.6030	

The Phase function in the Tertiary Scattering (1)

Table 1.  $\theta = 60$

$\theta_1$	$\theta_2$	$\theta_3$	0					30					60					90				
			0	1	2	3	4	0	1	2	3	4	0	1	2	3	4	0	1	2	3	4
30	0		equal to $\theta = 0$																			
	30		equal to $\theta = 0$																			
	60		equal to $\theta = 0$																			
	75		equal to $\theta = 0$																			
	90		equal to $\theta = 0$																			
	105		equal to $\theta = 0$																			
	120		equal to $\theta = 0$																			
	180		equal to $\theta = 0$																			
60	0		1.0000					1.0000														
	30		1.0000					1.0000														
	60		1.0000					1.0000														
	75		1.0000					1.0000														
	90		1.0000					1.0000														
	105		1.0000					1.0000														
	120		1.0000					1.0000														
	180		1.0000					1.0000														
75	0		1.0000	1.0000	1.0000	1.0000	1.0000	1.0011	1.0010	1.0000	1.0009	1.0008	1.0034	1.0030	1.0026	1.0023	1.0020	1.0045	1.0039	1.0034	1.0029	1.0024
	30		1.0011	1.0014	1.0016	1.0019	1.0022	1.0034	1.0039	1.0045	1.0051	1.0057	1.0045	1.0050	1.0057	1.0064	1.0071	1.0034	1.0038	1.0041	1.0045	1.0049
	60		1.0251	1.0265	1.0279	1.0294	1.0308	1.0335	1.0347	1.0358	1.0369	1.0380	1.0251	1.0255	1.0258	1.0260	1.0262	1.0084	1.0081	1.0079	1.0076	1.0073
	75		1.0469	1.0483	1.0496	1.0509	1.0521	1.0469	1.0474	1.0478	1.0483	1.0484	1.0251	1.0247	1.0243	1.0238	1.0233	1.0034	1.0030	1.0026	1.0023	1.0020
	90		1.0625	1.0631	1.0635	1.0639	1.0642	1.0469	1.0464	1.0458	1.0451	1.0444	1.0156	1.0148	1.0140	1.0132	1.0124	1.0000	1.0000	1.0001	1.0002	1.0003
	105		1.0625	1.0619	1.0612	1.0604	1.0596	1.0339	1.0323	1.0311	1.0299	1.0287	1.0045	1.0039	1.0034	1.0029	1.0024	1.0045	1.0051	1.0057	1.0064	1.0071
	120		1.0469	1.0455	1.0441	1.0426	1.0412	1.0156	1.0145	1.0135	1.0125	1.0115	1.0000	1.0000	1.0001	1.0002	1.0003	1.0156	1.0164	1.0172	1.0180	1.0187
	180		1.0084	1.0076	1.0069	1.0062	1.0055	1.0000	1.0000	1.0000	1.0001	1.0001	1.0156	1.0164	1.0172	1.0180	1.0187	1.0251	1.0247	1.0243	1.0238	1.0233
90	0		1.0000	1.0000	1.0000	1.0000	1.0000	1.0000	1.0000	1.0001	1.0002	1.0004	1.0000	1.0001	1.0002	1.0006	1.0010	1.0000	1.0001	1.0003	1.0007	1.0012
	30		1.0156	1.0176	1.0197	1.0220	1.0245	1.0469	1.0507	1.0547	1.0589	1.0632	1.0625	1.0663	1.0701	1.0740	1.0778	1.0469	1.0487	1.0505	1.0522	1.0538
	60		1.1406	1.1463	1.1520	1.1576	1.1632	1.1875	1.1912	1.1947	1.1979	1.2010	1.1334	1.1405	1.1402	1.1396	1.1388	1.0469	1.0450	1.0430	1.0409	1.0388
	75		1.2176	1.2216	1.2254	1.2289	1.2321	1.2176	1.2175	1.2171	1.2164	1.2154	1.1166	1.1136	1.1104	1.1071	1.1037	1.0156	1.0138	1.0120	1.0103	1.0088
	90		1.2500	1.2499	1.2494	1.2486	1.2476	1.1875	1.1837	1.1796	1.1754	1.1710	1.0625	1.0587	1.0550	1.0514	1.0478	1.0000	1.0001	1.0003	1.0007	1.0012
	105		1.2176	1.2135	1.2091	1.2046	1.1998	1.1166	1.1115	1.1063	1.1013	1.0962	1.0156	1.0135	1.0116	1.0098	1.0081	1.0156	1.0176	1.0196	1.0216	1.0237
	120		1.1406	1.1350	1.1293	1.1237	1.1181	1.0469	1.0432	1.0396	1.0363	1.0330	1.0000	1.0002	1.0002	1.0005	1.0009	1.0469	1.0487	1.0504	1.0522	1.0537
	180		1.0156	1.0138	1.0121	1.0106	1.0093	1.0000	1.0000	1.0001	1.0001	1.0003	1.0084	1.0156	1.0155	1.0153	1.0150	1.0469	1.0450	1.0430	1.0409	1.0388



105	0	1.0000	1.0000	1.0001	1.0001	1.0003	1.0084	1.0102	1.0102	1.0142	1.0166	1.0252	1.0292	1.0335	1.0380	1.0429	1.0333	1.0380	1.0427	1.0478	1.0527
	30	1.0625	1.0687	1.0751	1.0819	1.0893	1.1876	1.1979	1.2085	1.2192	1.2301	1.2501	1.2587	1.2672	1.2755	1.2835	1.1876	1.1901	1.1923	1.1943	1.1958
	60	1.3499	1.3601	1.3701	1.3798	1.3892	1.4665	1.4707	1.4742	1.4771	1.4795	1.3499	1.3459	1.3415	1.3365	1.3311	1.1167	1.1106	1.1046	1.0986	1.0926
	75	1.4665	1.4706	1.4742	1.4772	1.4801	1.4665	1.4619	1.4567	1.4511	1.4448	1.2500	1.2413	1.2323	1.2233	1.2142	1.0335	1.0293	1.0253	1.0216	1.0181
	90	1.4665	1.4620	1.4567	1.4510	1.4449	1.3499	1.3396	1.3289	1.3181	1.3072	1.1167	1.1086	1.1008	1.0932	1.0859	1.0000	1.0002	1.0006	1.0012	1.0022
	105	1.3499	1.3395	1.3289	1.3182	1.3072	1.1875	1.1773	1.1673	1.1575	1.1480	1.0252	1.0215	1.0182	1.0151	1.0125	1.0252	1.0279	1.0308	1.0336	1.0365
	120	1.1876	1.1774	1.1673	1.1575	1.1479	1.0625	1.0567	1.0512	1.0462	1.0414	1.0000	1.0001	1.0003	1.0006	1.0011	1.0625	1.0640	1.0653	1.0664	1.0673
	150	1.0084	1.0069	1.0056	1.0045	1.0035	1.0000	1.0001	1.0000	1.0001	1.0001	1.0084	1.0078	1.0072	1.0064	1.0057	1.0252	1.0225	1.0197	1.0172	1.0147
	180											1.0252	1.0292	1.0335	1.0380	1.0429	1.0333	1.0380	1.0427	1.0478	1.0527
120	0	equal to $\theta = 0$																			
	30																				
	60																				
	75																				
	90																				
	180																				
150	0	1.0000	1.0002	1.0009	1.0026	1.0040	1.1875	1.2028	1.2189	1.2354	1.2526	1.5625	1.5851	1.6078	1.6303	1.6526	1.7500	1.7648	1.7787	1.7917	1.8040
	30	1.2500	1.2651	1.2805	1.2958	1.3112	1.7500	1.7647	1.7787	1.7917	1.8039	2.0000	1.9995	1.9976	1.9945	1.9904	1.7500	1.7346	1.7184	1.7015	1.6838
	60	1.5625	1.5620	1.5608	1.5585	1.5552	1.7500	1.7346	1.7183	1.7014	1.6839	1.5625	1.5400	1.5173	1.4947	1.4723	1.1875	1.1726	1.1585	1.1450	1.1321
	75	1.4665	1.4544	1.4416	1.4284	1.4149	1.4665	1.4459	1.4253	1.4052	1.3848	1.2500	1.2329	1.2163	1.2005	1.1853	1.0335	1.0282	1.0235	1.0193	1.0157
	90	1.2500	1.2350	1.2200	1.2055	1.1913	1.1875	1.1726	1.1584	1.1449	1.1320	1.0625	1.0552	1.0485	1.0425	1.0370	1.0000	1.0001	1.0003	1.0005	1.0010
	105	1.0625	1.0540	1.0462	1.0391	1.0324	1.0335	1.0283	1.0234	1.0193	1.0156	1.0045	1.0034	1.0026	1.0018	1.0013	1.0045	1.0045	1.0043	1.0041	1.0039
	120	1.0000	1.0004	1.0009	1.0020	1.0034	1.0000	1.0001	1.0003	1.0005	1.0010	1.0000	1.0000	1.0000	1.0000	1.0001	1.0000	1.0001	1.0003	1.0008	1.0014
	150	1.0625	1.0624	1.0618	1.0611	1.0601	1.0000	1.0001	1.0003	1.0008	1.0016	1.0625	1.0704	1.0788	1.0879	1.0978	1.1875	1.2028	1.2188	1.2354	1.2526
	180											1.5625	1.5851	1.6078	1.6303	1.6526	1.7500	1.7648	1.7787	1.7917	1.8040

The Phase function in the Tertiary Scattering (1)

Table 1. A = 0  $\theta = 90$

$\theta_1$	$\theta_2$	$\theta_3$	h					30					60					90				
			0	1	2	3	4	0	1	2	3	4	0	1	2	3	4	0	1	2	3	4
30	0	0	1.0000	1.0002	1.0007	1.0014	1.0026	1.1406	1.1462	1.1515	1.1565	1.1612	1.4219	1.4214	1.4207	1.4189	1.4164	1.5625	1.5510	1.5387	1.5265	1.5129
	30	0	1.0469	1.0467	1.0464	1.0458	1.0451	1.1406	1.1349	1.1289	1.1227	1.1165	1.1875	1.1764	1.1650	1.1540	1.1435	1.1406	1.1295	1.1188	1.1087	1.0990
	60	0	1.0000	1.0002	1.0007	1.0017	1.0029	1.0000	1.0002	1.0009	1.0021	1.0036	1.0000	1.0003	1.0007	1.0015	1.0026	1.0000	1.0001	1.0002	1.0004	1.0008
	75	0	1.0469	1.0537	1.0608	1.0685	1.0766	1.0469	1.0526	1.0587	1.0647	1.0711	1.0251	1.0275	1.0298	1.0321	1.0344	1.0034	1.0033	1.0032	1.0031	1.0029
	90	0	1.1875	1.1989	1.2104	1.2220	1.2335	1.1406	1.1462	1.1514	1.1565	1.1613	1.0469	1.0468	1.0464	1.0458	1.0451	1.0000	1.0001	1.0002	1.0012	1.0012
	105	0	1.3499	1.3586	1.3668	1.3746	1.3821	1.1875	1.1873	1.1866	1.1855	1.1839	1.0251	1.0227	1.0203	1.0179	1.0155	1.0252	1.0297	1.4344	1.0396	1.0453
	120	0	1.4219	1.4216	1.4205	1.4189	1.4164	1.1406	1.1349	1.1289	1.1227	1.1165	1.0000	1.0002	1.0007	1.0016	1.0030	1.1406	1.1521	1.1642	1.1766	1.1895
150	0	1.1875	1.1762	1.1651	1.1541	1.1435	1.0000	1.0002	1.0009	1.0020	1.0037	1.1875	1.1988	1.2104	1.2219	1.2334	1.5625	1.5735	1.5841	1.5938	1.6029	
60	0	0																				
	30	0																				
	60	0																				
	75	0																				
	90	0																				
	105	0																				
	150	0																				
75	0	0																				
	30	0																				
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	150	0																				
90	0	0																				
	30	0																				
	60	0																				
	75	0																				
	90	0																				
	105	0																				
	150	0																				

105	0	1.0000	1.0000	1.0000	1.0000	1.0000	1.0011	1.0014	1.0016	1.0019	1.0022	1.0034	1.0039	1.0045	1.0051	1.0057	1.0045	1.0051	1.0057	1.0064	1.0071
	30	1.0084	1.0092	1.0101	1.0110	1.0120	1.0251	1.0265	1.0279	1.0294	1.0308	1.0335	1.0347	1.0358	1.0369	1.0011	1.0251	1.0255	1.0258	1.0260	1.0262
	60	1.0469	1.0483	1.0496	1.0509	1.0522	1.0625	1.0631	1.0635	1.0639	1.0642	1.0469	1.0464	1.0458	1.0451	1.0444	1.0156	1.0148	1.0140	1.0132	1.0124
	75	1.0625	1.0631	1.0635	1.0639	1.0643	1.0625	1.0619	1.0612	1.0604	1.0596	1.0355	1.0323	1.0311	1.0299	1.0287	1.0045	1.0039	1.0034	1.0029	1.0024
	90	1.0625	1.0619	1.0612	1.0604	1.0596	1.0469	1.0455	1.0441	1.0426	1.0412	1.0156	1.0145	1.0135	1.0125	1.0115	1.0000	1.0000	1.0001	1.0002	1.0003
	105	1.0469	1.0455	1.0441	1.0426	1.0412	1.0251	1.0238	1.0224	1.0211	1.0198	1.0034	1.0029	1.0024	1.0020	1.0017	1.0034	1.0037	1.0041	1.0045	1.0049
	120	1.0251	1.0238	1.0224	1.0211	1.0198	1.0084	1.0076	1.0069	1.0062	1.0055	1.0000	1.0000	1.0000	1.0001	1.0001	1.0084	1.0086	1.0088	1.0089	1.0090
150	1.0011	1.0009	1.0007	1.0006	1.0005	1.0000	1.0000	1.0000	1.0000	1.0000	1.0011	1.0010	1.0010	1.0009	1.0008	1.0034	1.0030	1.0026	1.0023	1.0020	
											1.0034	1.0039	1.0045	1.0051	1.0057	1.0045	1.0051	1.0057	1.0064	1.0071	
120	0	1.0000	1.0000	1.0001	1.0002	1.0004	1.0156	1.0176	1.0197	1.0220	1.0245	1.0469	1.0507	1.0547	1.0589	1.0632	1.0625	1.0663	1.0701	1.0740	1.0078
	30	1.0469	1.0507	1.0547	1.0589	1.0632	1.1406	1.1463	1.1520	1.1576	1.1632	1.1875	1.1912	1.1947	1.1979	1.2012	1.1406	1.1405	1.1402	1.1396	1.1388
	60	1.1875	1.1912	1.1947	1.1979	1.2010	1.2500	1.2499	1.2494	1.2486	1.2476	1.1875	1.1837	1.1796	1.1754	1.1710	1.0625	1.0587	1.0550	1.0514	1.0478
	75	1.2176	1.2175	1.2171	1.2164	1.2154	1.2176	1.2135	1.2091	1.2046	1.1998	1.1166	1.1115	1.1063	1.1013	1.0962	1.0156	1.0135	1.0116	1.0098	1.0081
	90	1.1875	1.1837	1.1796	1.1754	1.1710	1.1406	1.1350	1.1293	1.1237	1.1181	1.0469	1.0432	1.0396	1.0363	1.0330	1.0000	1.0001	1.0002	1.0005	1.0009
	105	1.1166	1.1115	1.1063	1.1013	1.0962	1.0625	1.0582	1.0541	1.0501	1.0463	1.0084	1.0071	1.0059	1.0048	1.0039	1.0084	1.0092	1.0100	1.0107	1.0114
	120	1.0469	1.0432	1.0396	1.0362	1.0330	1.0156	1.0138	1.0121	1.0106	1.0093	1.0000	1.0000	1.0001	1.0001	1.0003	1.0156	1.0156	1.0155	1.0153	1.0150
150	1.0000	1.0000	1.0001	1.0001	1.0003	1.0000	1.0000	1.0000	1.0000	1.0001	1.0000	1.0000	1.0001	1.0002	1.0004	1.0000	1.0001	1.0002	1.0004	1.0010	
											1.0469	1.0507	1.0547	1.0589	1.0632	1.0625	1.0663	1.0701	1.0740	1.0078	
150	0	equal to $\theta = 30$																			
	30																				
	60																				
	75																				
	90																				
	105																				
	120																				
150																					